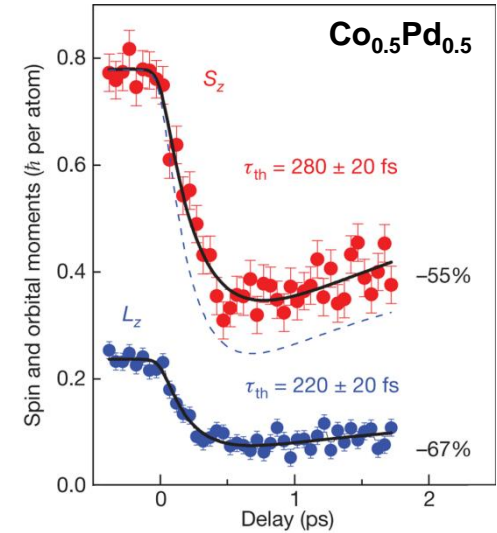
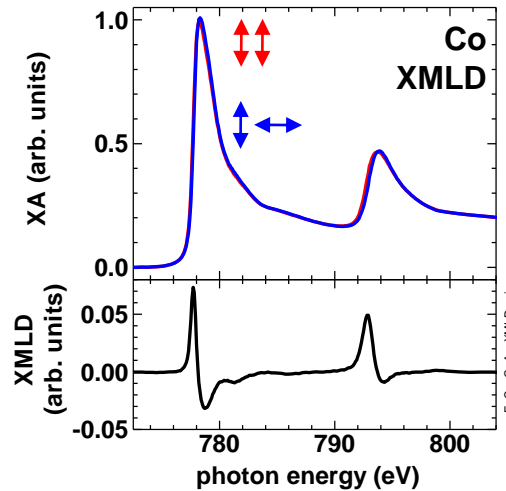
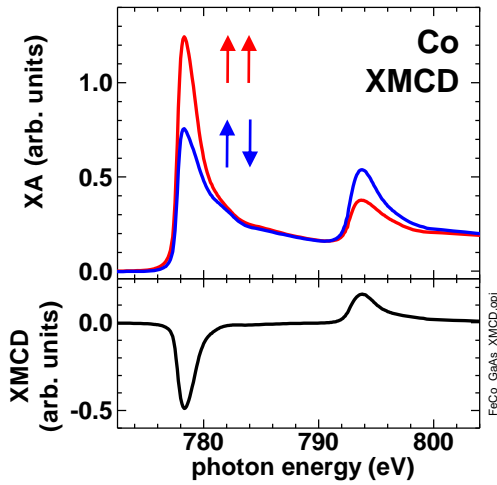
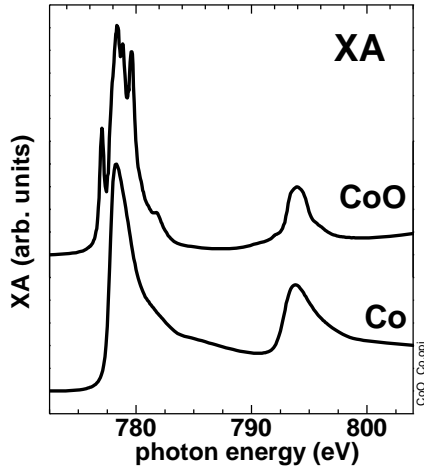
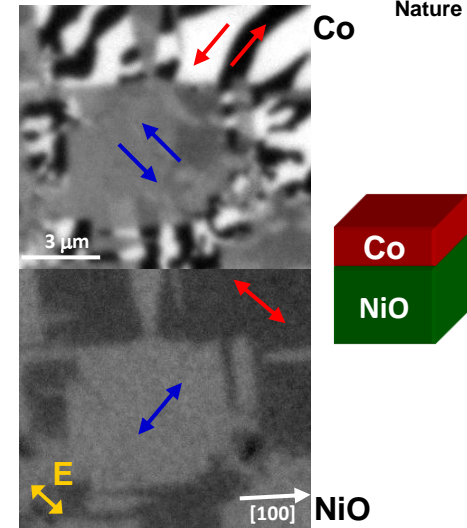


Elke Arenholz, Advanced Light Source

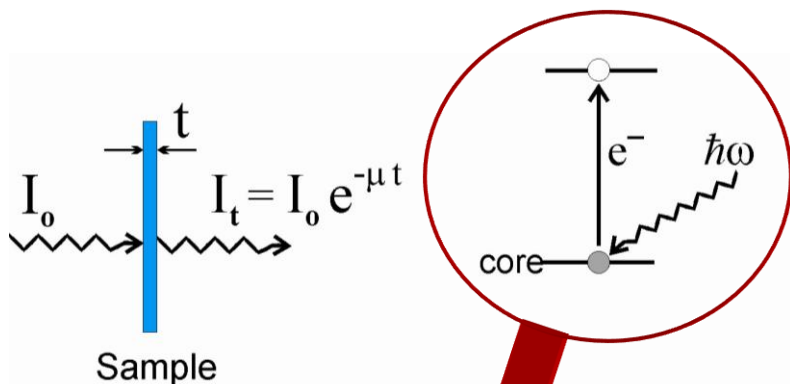
- + X-ray absorption, XA
- + X-ray magnetic circular dichroism, XMCD
- + X-ray magnetic linear dichroism, XMLD
- + X-ray magnetic microscopy
- + Magnetization Dynamics



C. Boeglin *et al.*,  
Nature **465**, 458 (2011)

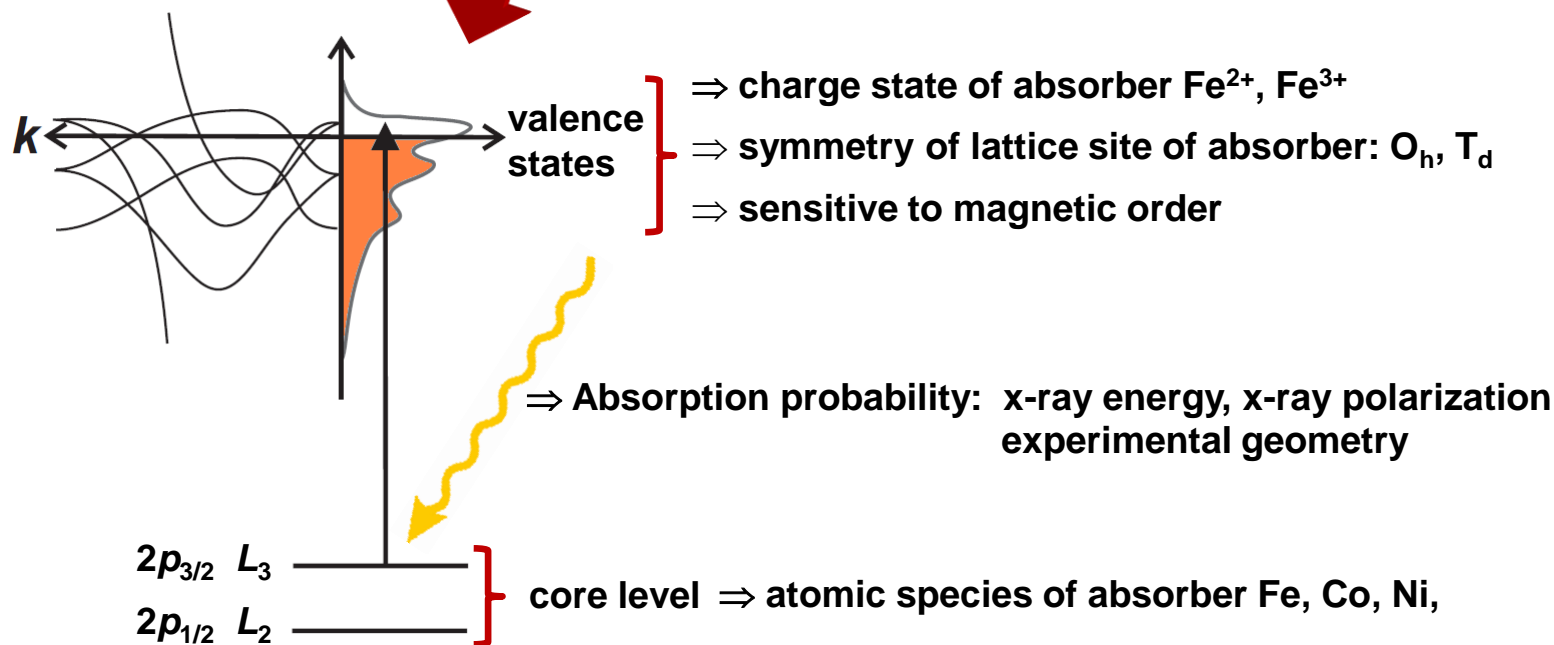


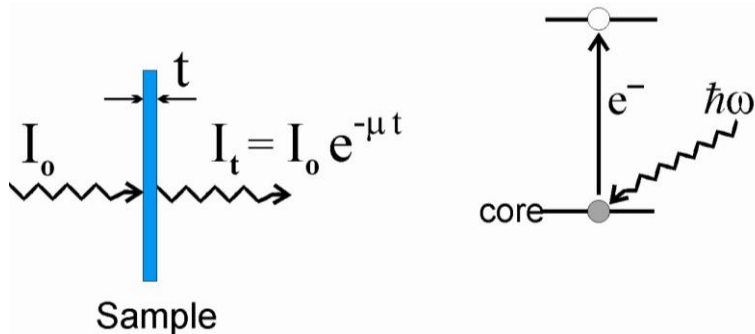
E. Arenholz *et al.*,  
Appl. Phys. Lett. **93**, 162506 (2008)



### Experimental Concept:

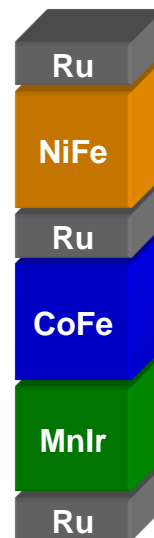
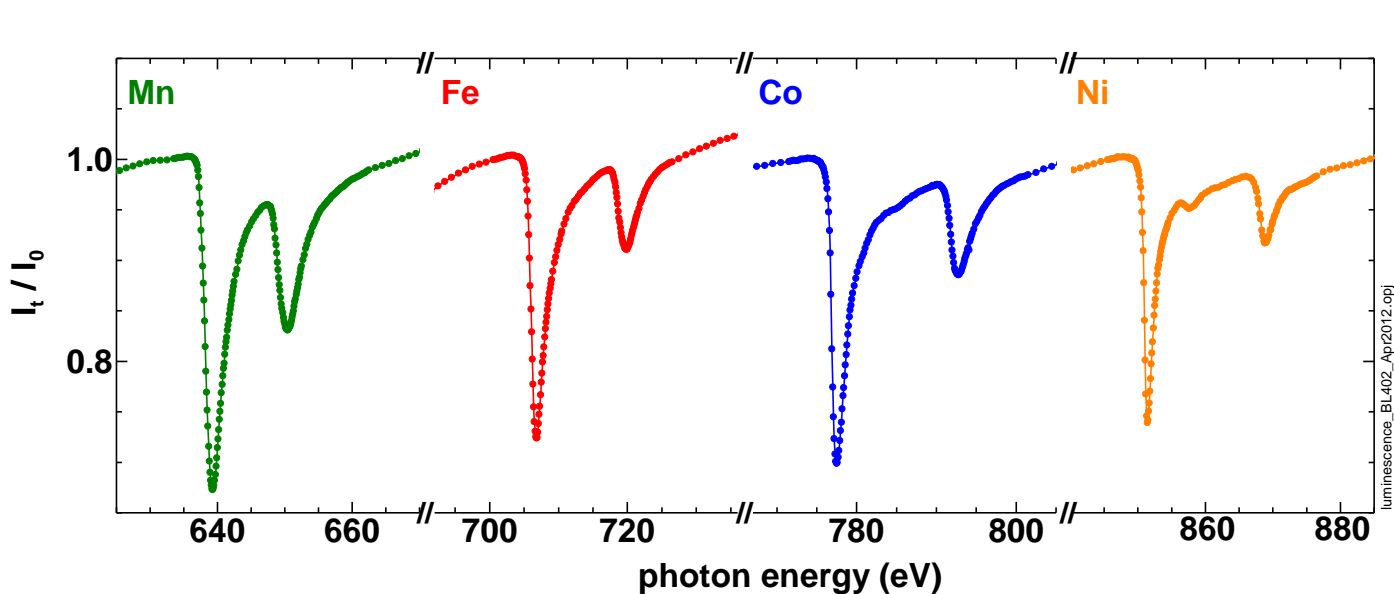
Monitor the reduction in x-ray flux transmitted through sample as function of x-ray photon energy



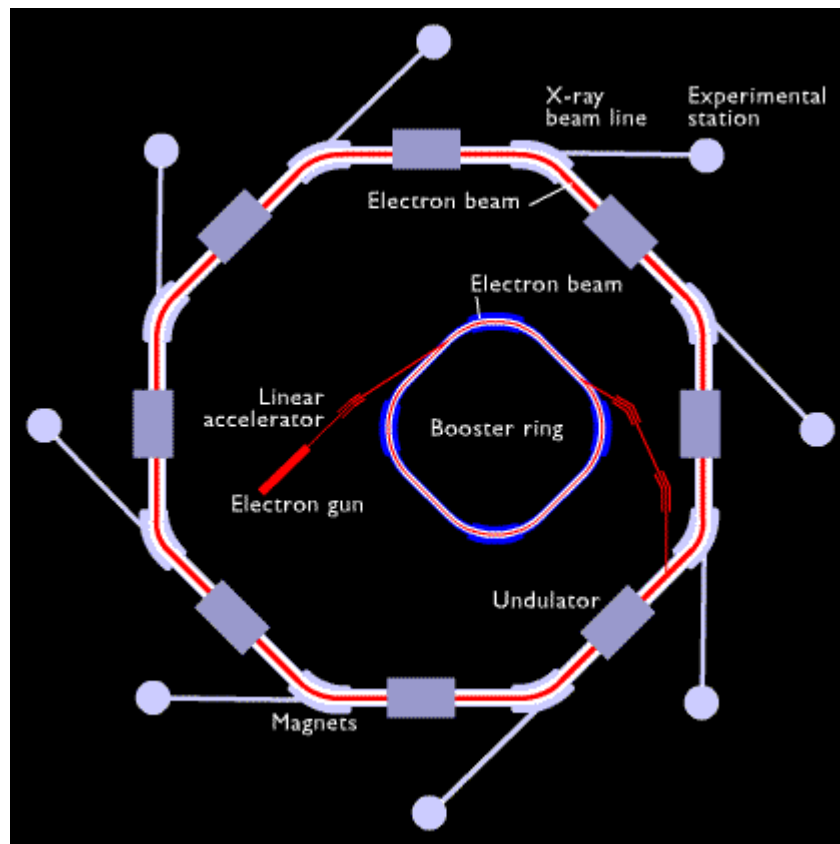


### Experimental Concept:

Monitor the reduction in x-ray flux transmitted through sample as function of x-ray photon energy



- + Change velocity of electrons near the speed of light
- ⇒ Emission of wide wavelength range of electromagnetic spectrum (IR, UV, soft x-rays, hard x-rays)
- ⇒ Tunable source in the soft x ray range for x-ray absorption spectroscopy

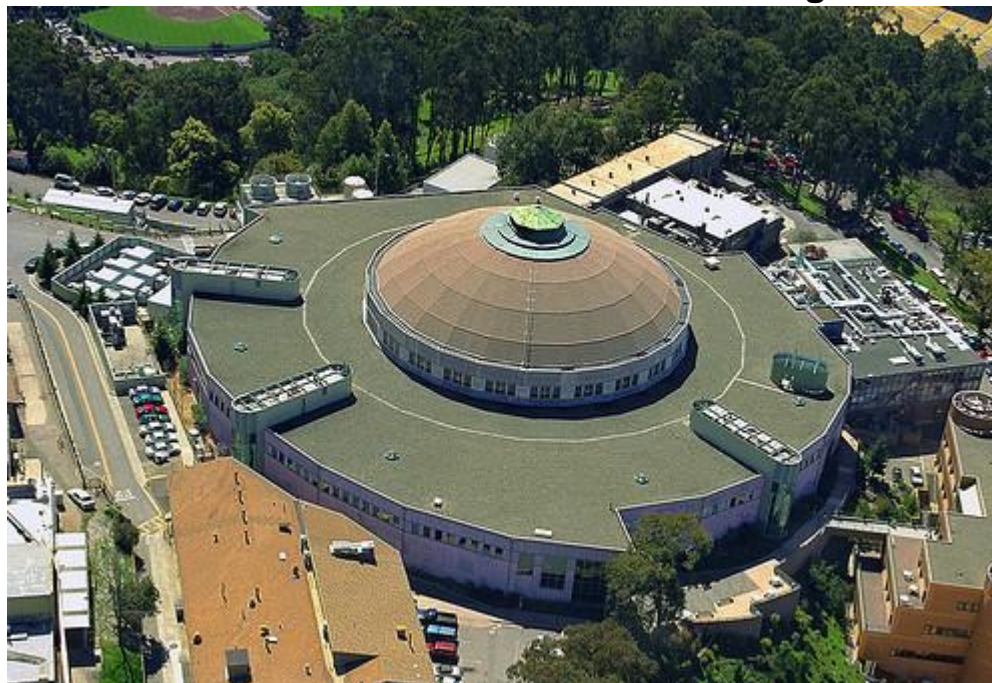


<http://www.ph.surrey.ac.uk/>



- + Change velocity of electrons near the speed of light
- ⇒ Emission of wide wavelength range of electromagnetic spectrum (IR, UV, soft x-rays, hard x-rays)
- ⇒ Tunable source in the soft x-ray range for x-ray absorption spectroscopy

Advanced Light Source



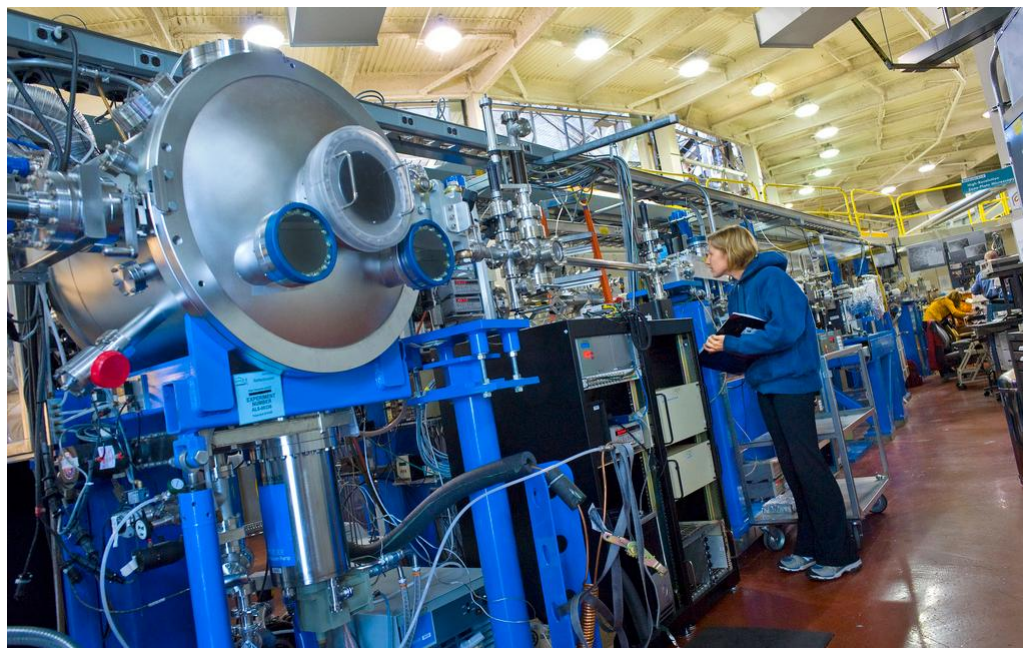


## Elliptically Polarized Undulator

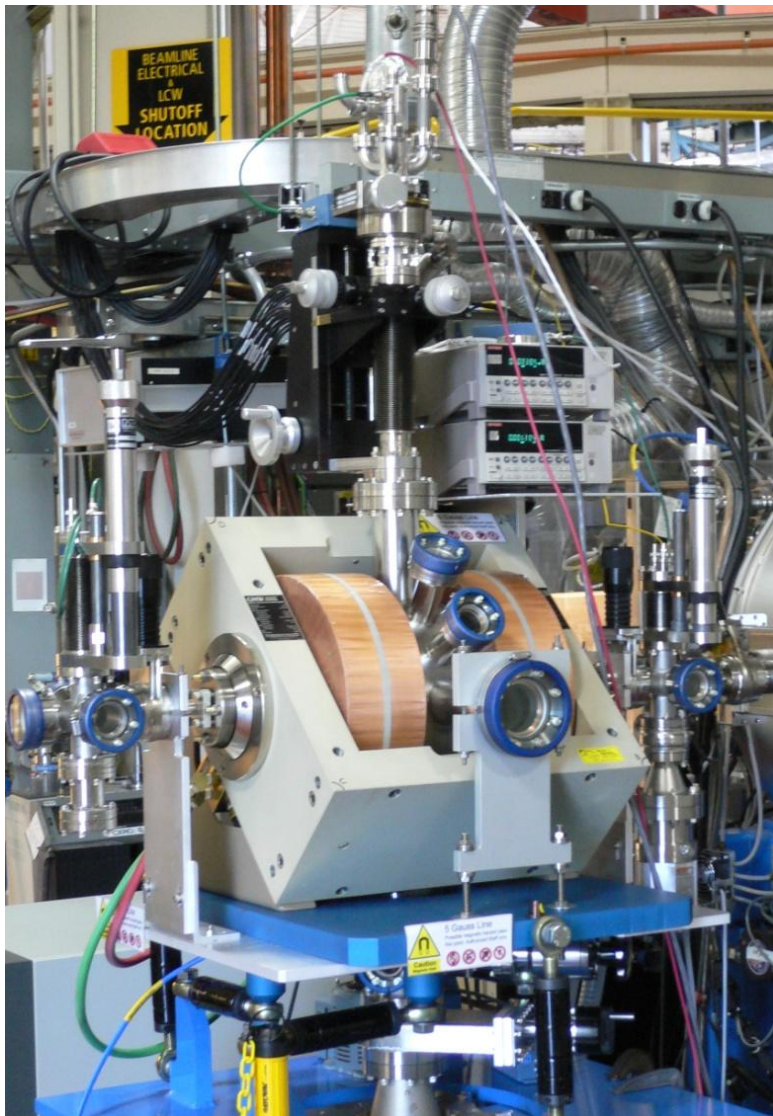


- + Beamlines/Monochromators provide photons with well defined characteristics:
  - tunable energy/wavelength
  - band width
  - fixed polarization: (variable) linear, circular, ...

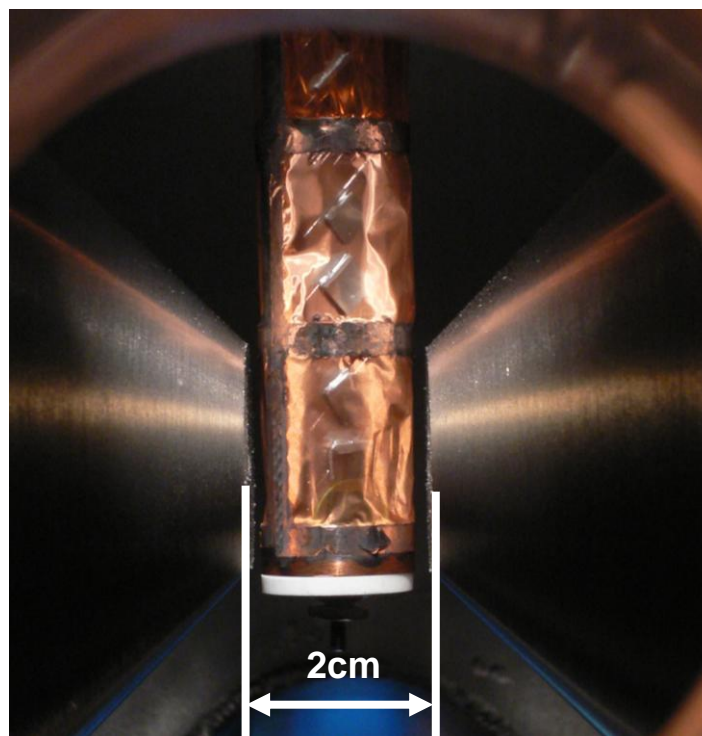
BL6.3.1



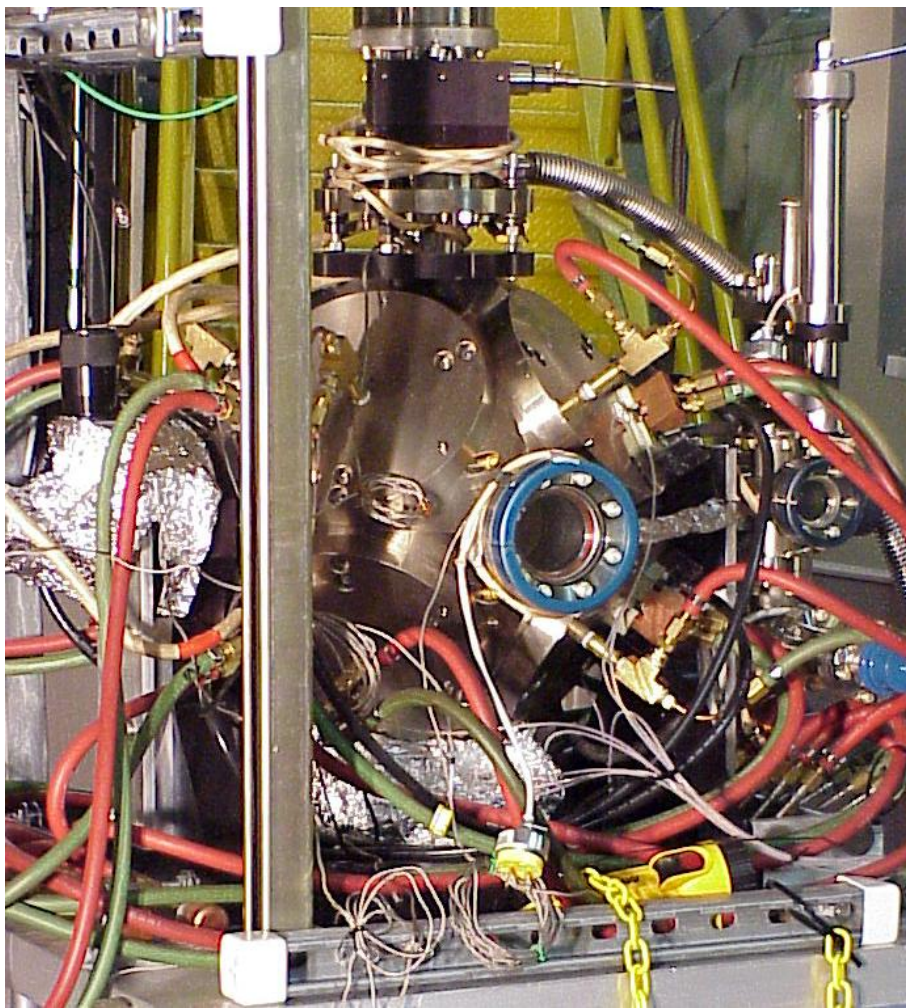




- + Endstations provide well defined environments for the interaction of samples and photons:
  - precisely defined experimental geometries ,  
i.e. angle of x ray beam to sample
  - sample temperature
  - external magnetic and electric fields

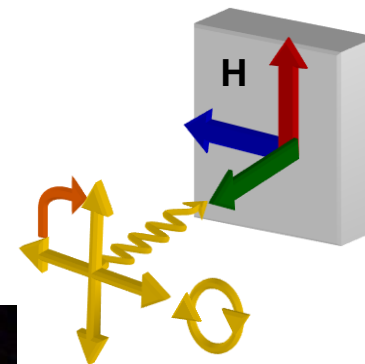
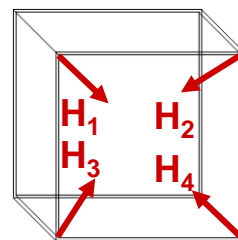


2T magnet at  
ALS BL6.3.1



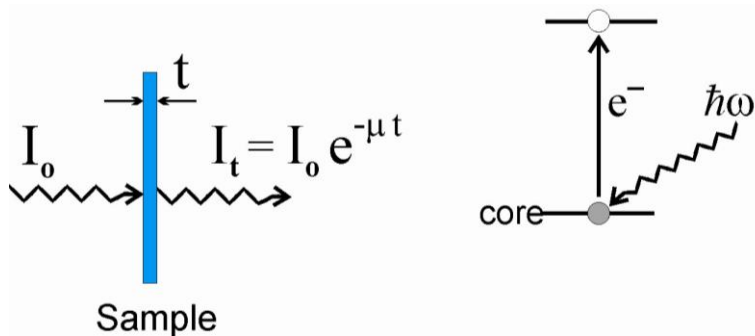
Vector magnet at ALS BL4.0.2

- + Magnetic fields in arbitrary directions obtained through superposition of fields generated by 4 dipole pairs in octahedral configuration.



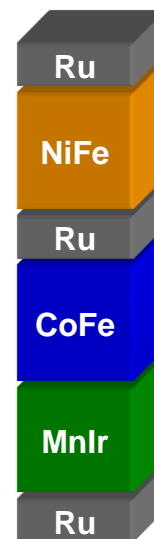
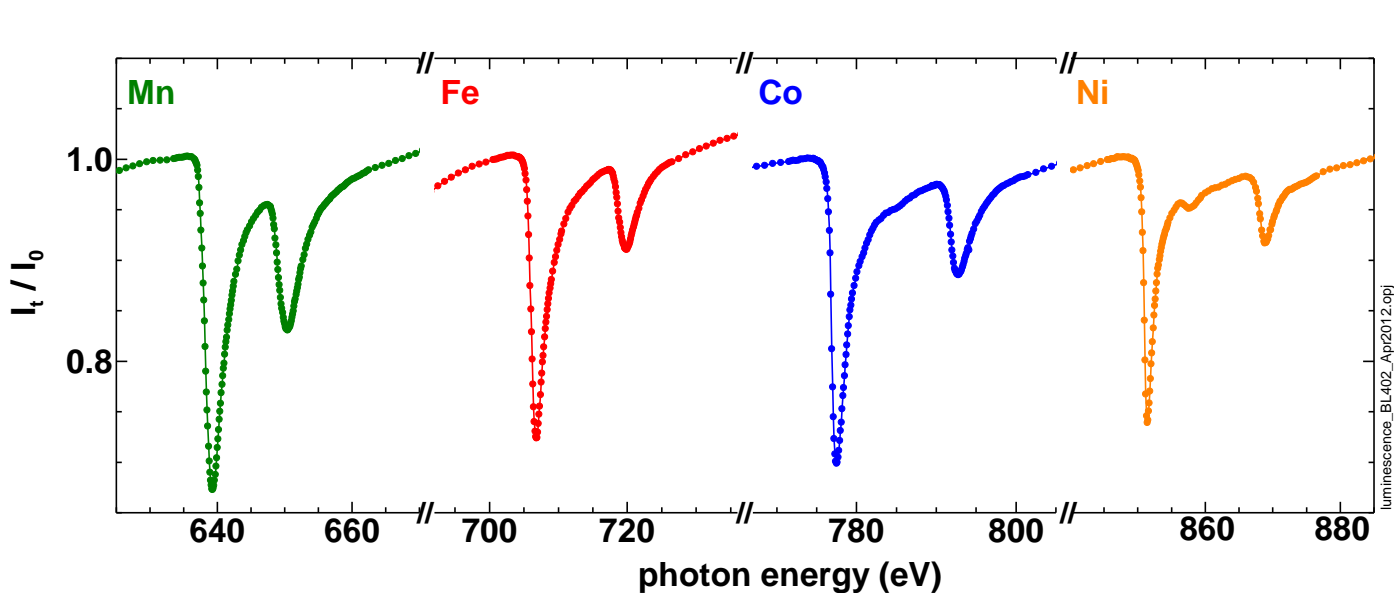
2cm

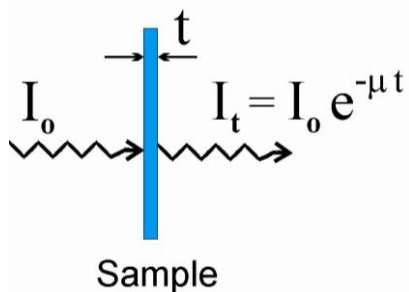




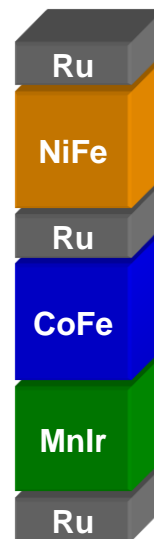
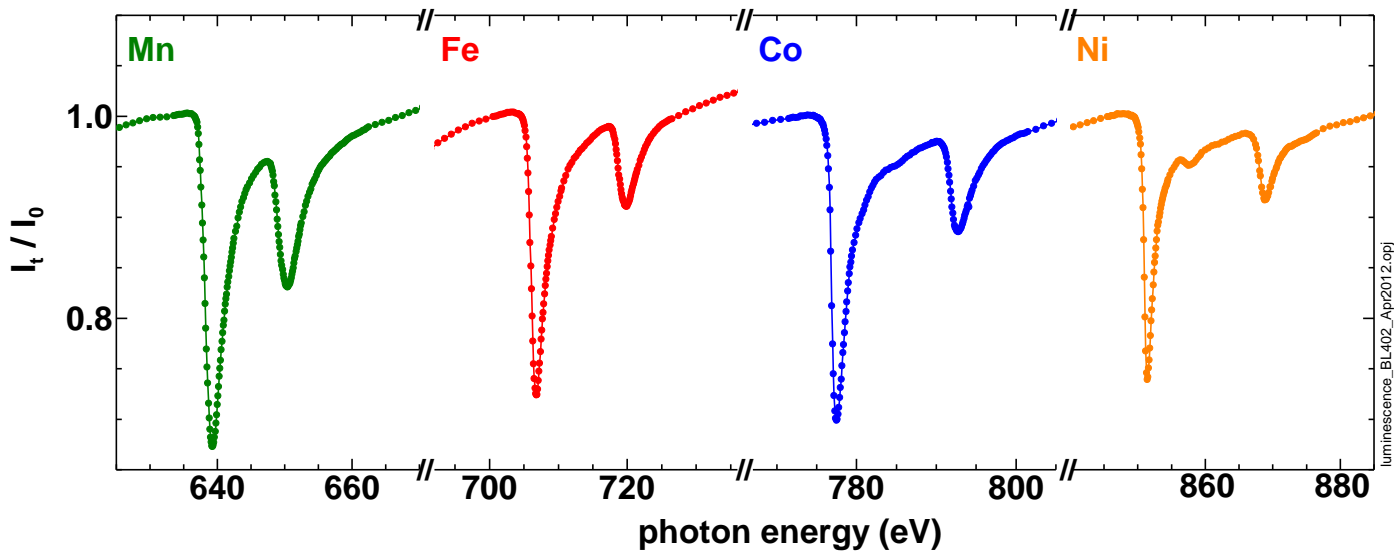
### Experimental Concept:

Monitor the reduction in x-ray flux transmitted through sample as function of x-ray photon energy

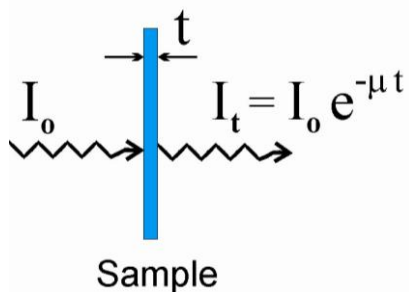




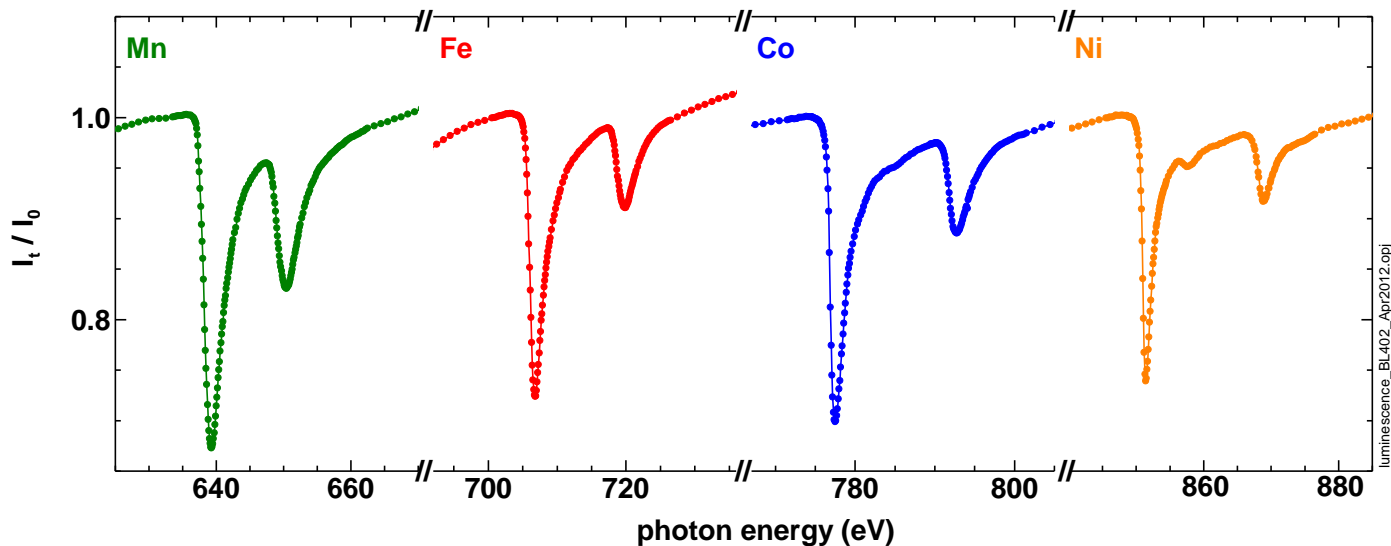
Element	10eV below $L_3$ $\mu$ [ $\text{nm}^{-1}$ ]	at $L_3$ $\mu$ [ $\text{nm}^{-1}$ ]	40 eV above $L_3$ $\mu$ [ $\text{nm}^{-1}$ ]
Fe	$1.8 \times 10^{-3}$	$6.0 \times 10^{-2}$	$1.2 \times 10^{-2}$
Co	$1.8 \times 10^{-3}$	$5.8 \times 10^{-2}$	$1.2 \times 10^{-2}$
Ni	$1.6 \times 10^{-3}$	$4.2 \times 10^{-2}$	$1.2 \times 10^{-2}$





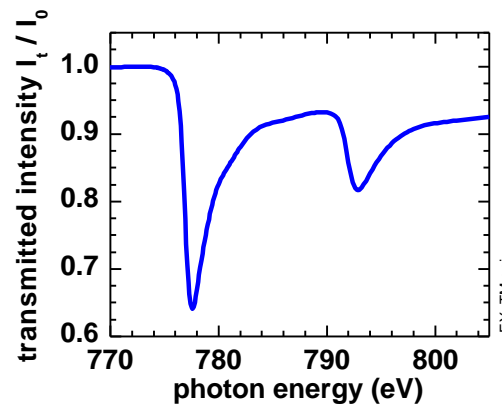
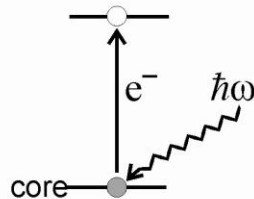
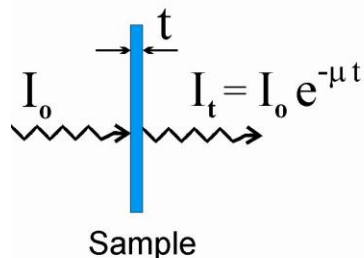


Element	10eV below $L_3$ $1/\mu$ [nm]	at $L_3$ $1/\mu$ [nm <sup>-1</sup> ]	40 eV above $L_3$ $1/\mu$ [nm <sup>-1</sup> ]
Fe	550	17	85
Co	550	17	85
Ni	625	24	85



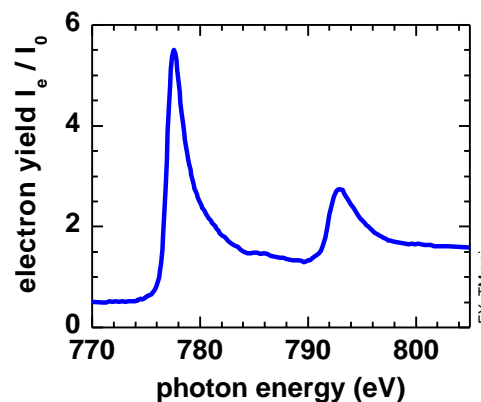
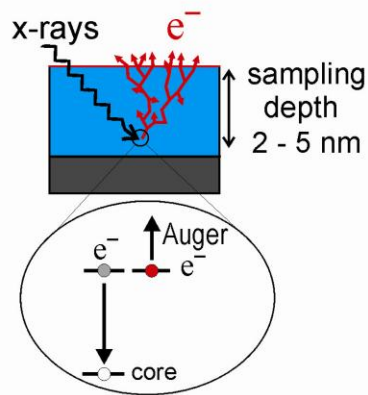
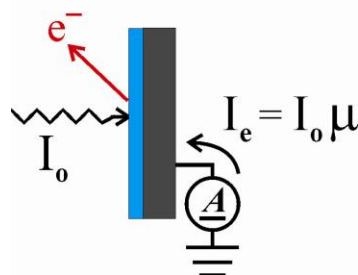
**~10-20 nm layer thick films supported by substrates transparent to soft x-rays**

## Transmission



photons  
absorbed

## Electron Yield



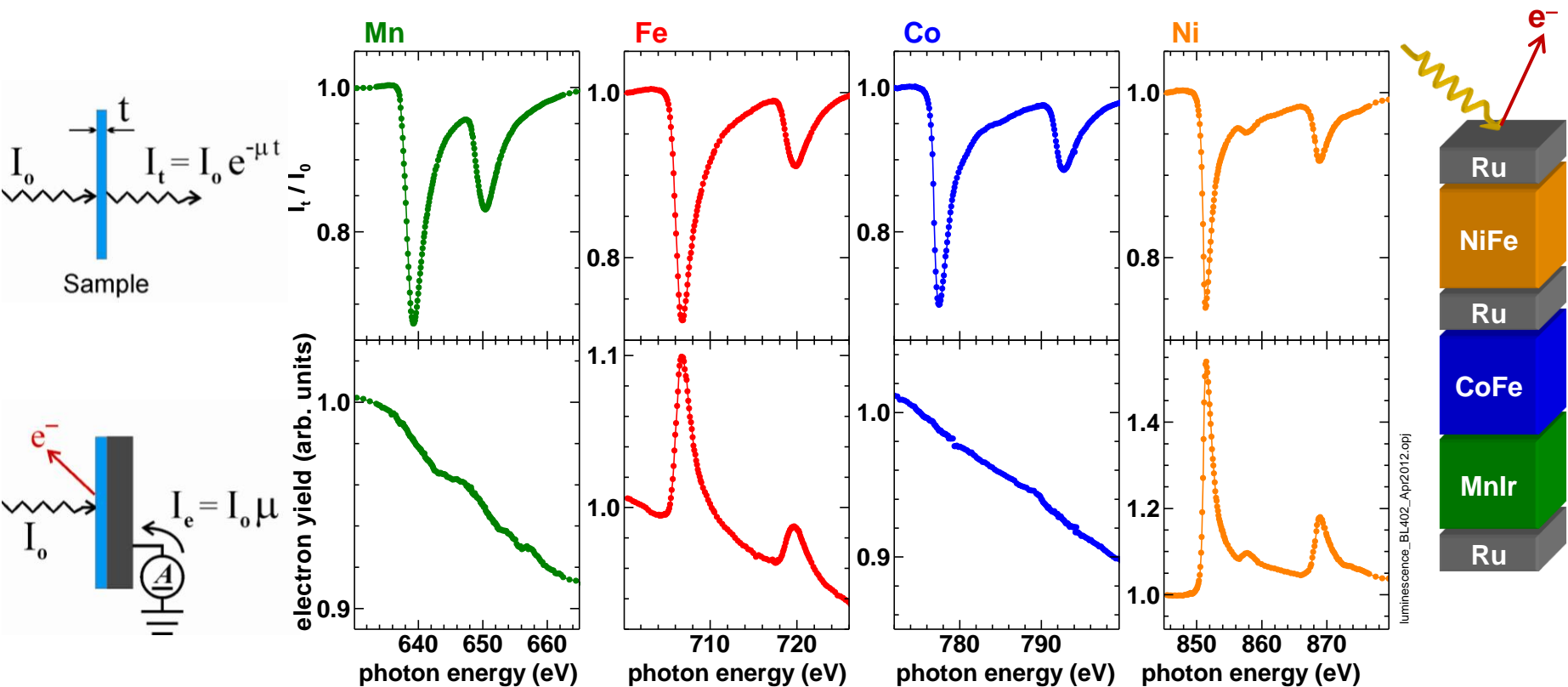
electrons  
generated

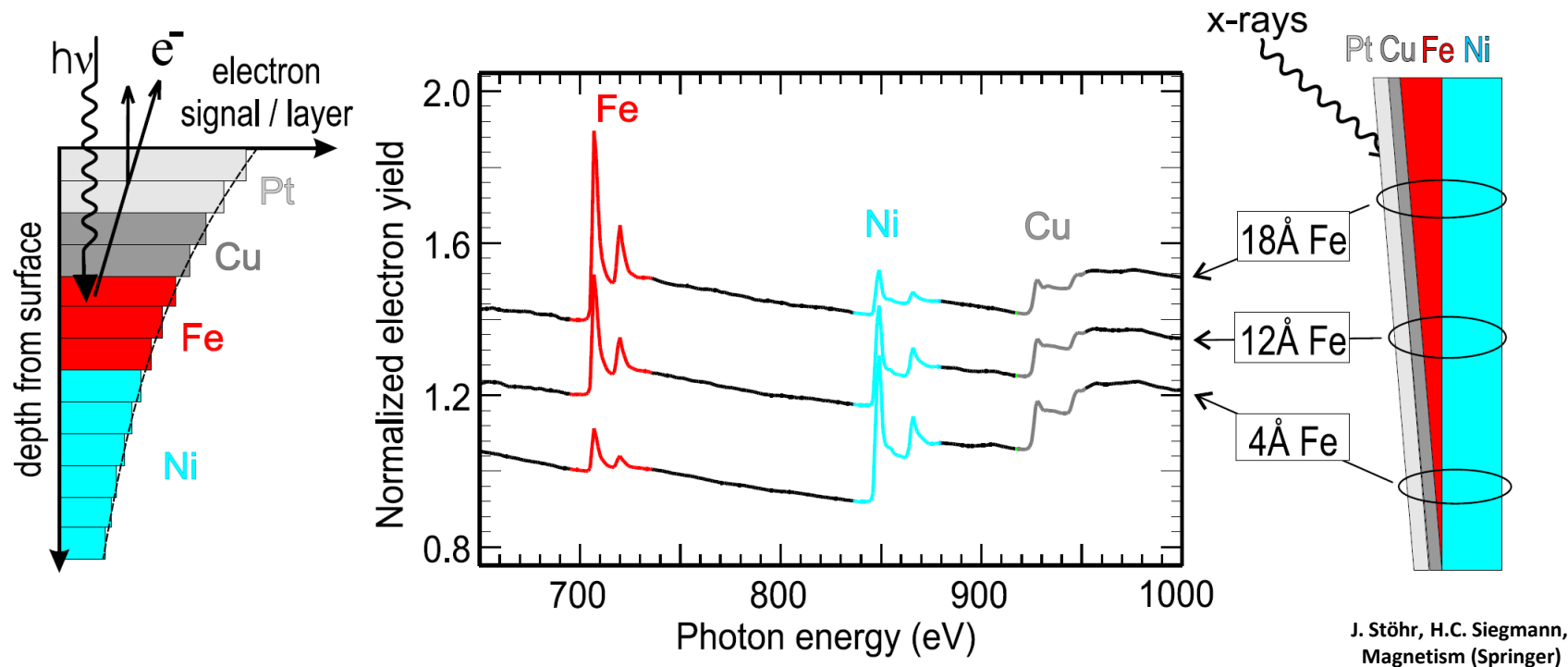
J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)

**Electron yield:**

- + Absorbed photons create core holes filled predominantly by Auger electron emission
- + Auger electrons create low-energy secondary electron cascade through inelastic scattering
- + Emitted electrons  $\propto$  probability of Auger electron creating  $\propto$  absorption probability



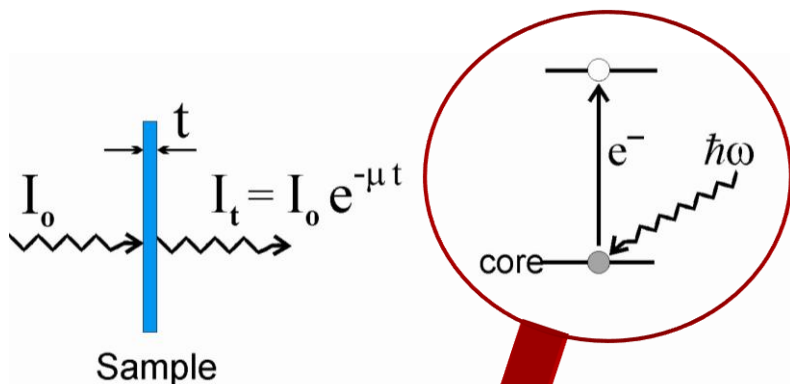




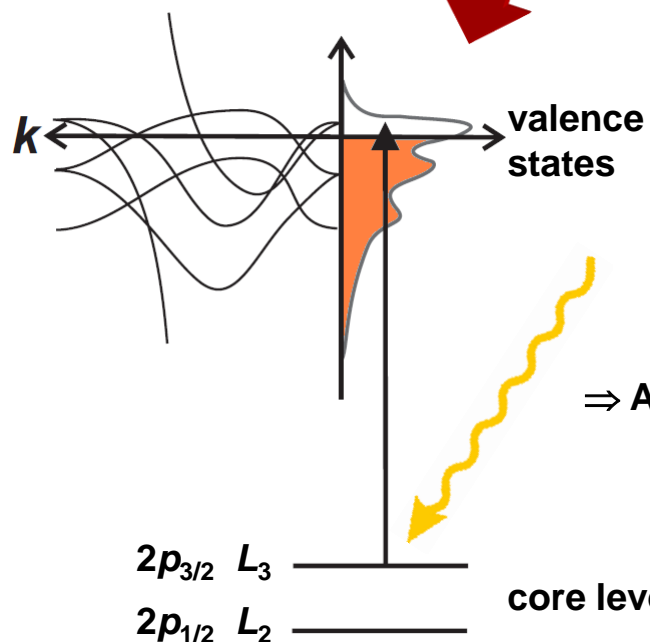
+ Electron sample depth: 2-5 nm in Fe, Co, Ni

⇒ 60% of the electron yield originates from the topmost 2-5 nm



**Experimental Concept:**

Monitor the reduction in x-ray flux transmitted through sample as function of x-ray photon energy



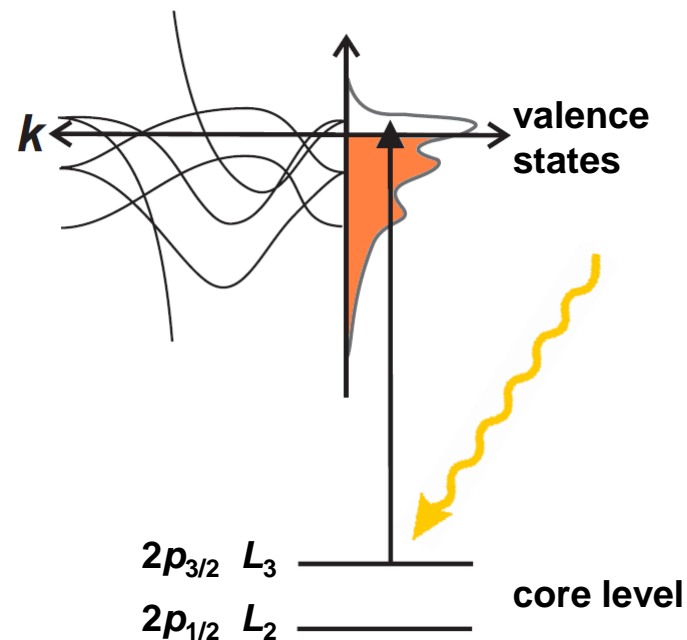
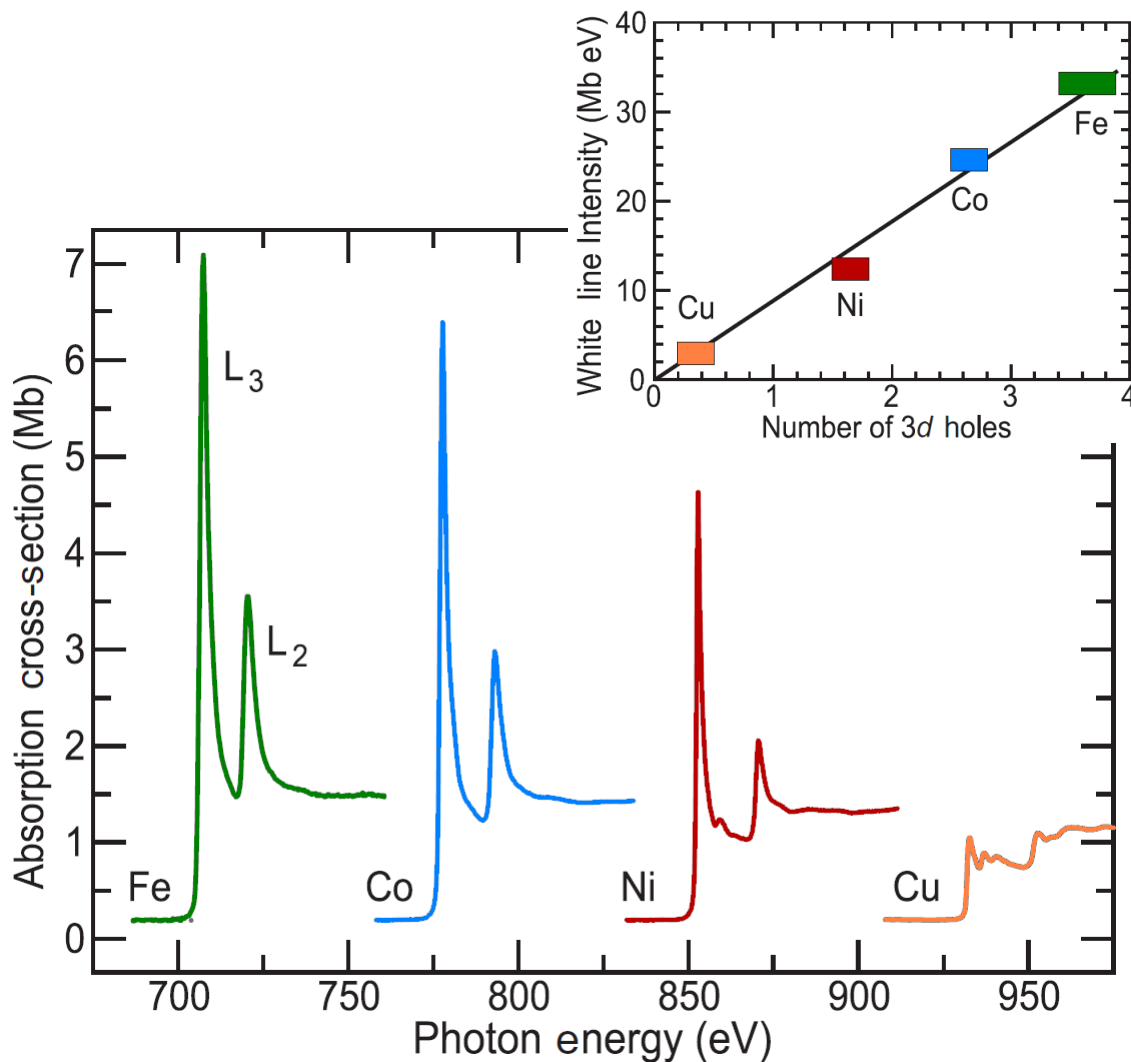
⇒ charge state of absorber  $\text{Fe}^{2+}$ ,  $\text{Fe}^{3+}$

⇒ symmetry of lattice site of absorber:  $\text{O}_h$ ,  $\text{T}_d$

⇒ sensitive to magnetic order

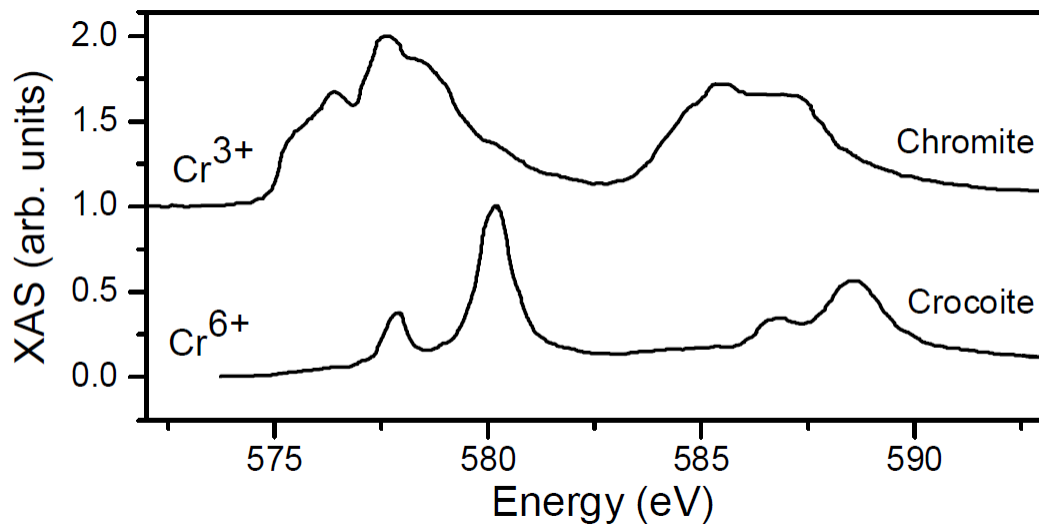
⇒ Absorption probability: x-ray energy, x-ray polarization  
experimental geometry

core level ⇒ atomic species of absorber Fe, Co, Ni,



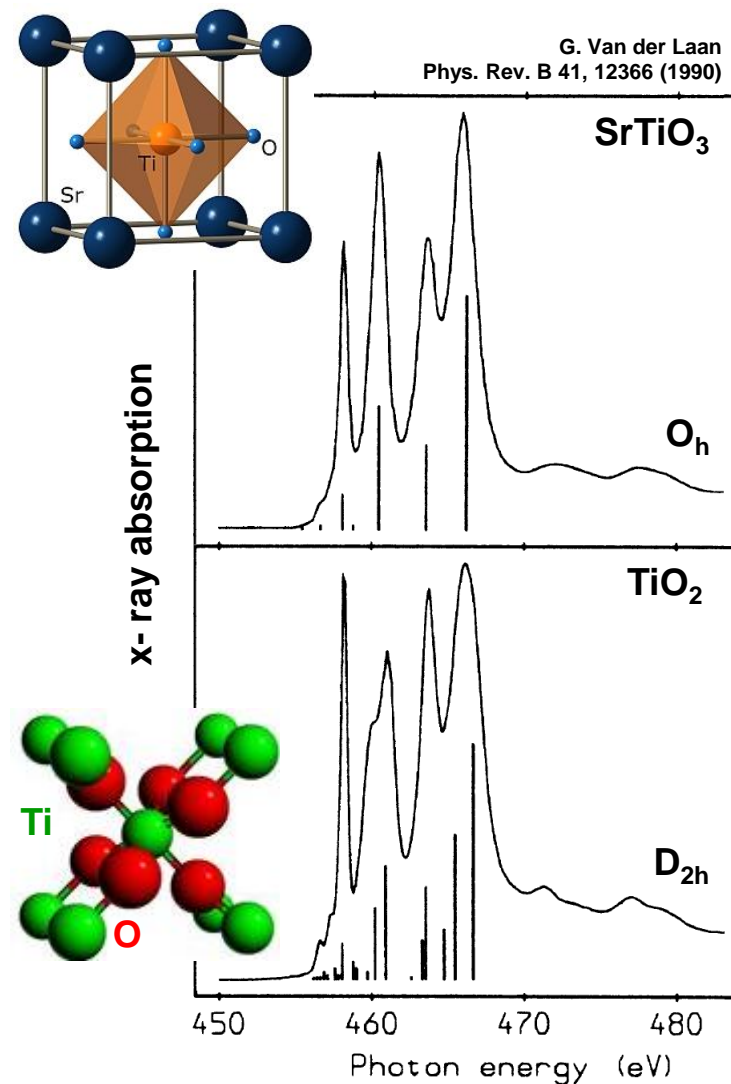
The intensity of the  $L_{3,2}$  resonances is proportional to the number of  $d$  states above the Fermi level, i.e. the number of holes in the  $d$  band.

### Influence of the charge state of the absorber



N. Telling *et al.*,  
*Appl. Phys. Lett.* **95**, 163701 (2009)

### Influence of lattice site symmetry at the absorber



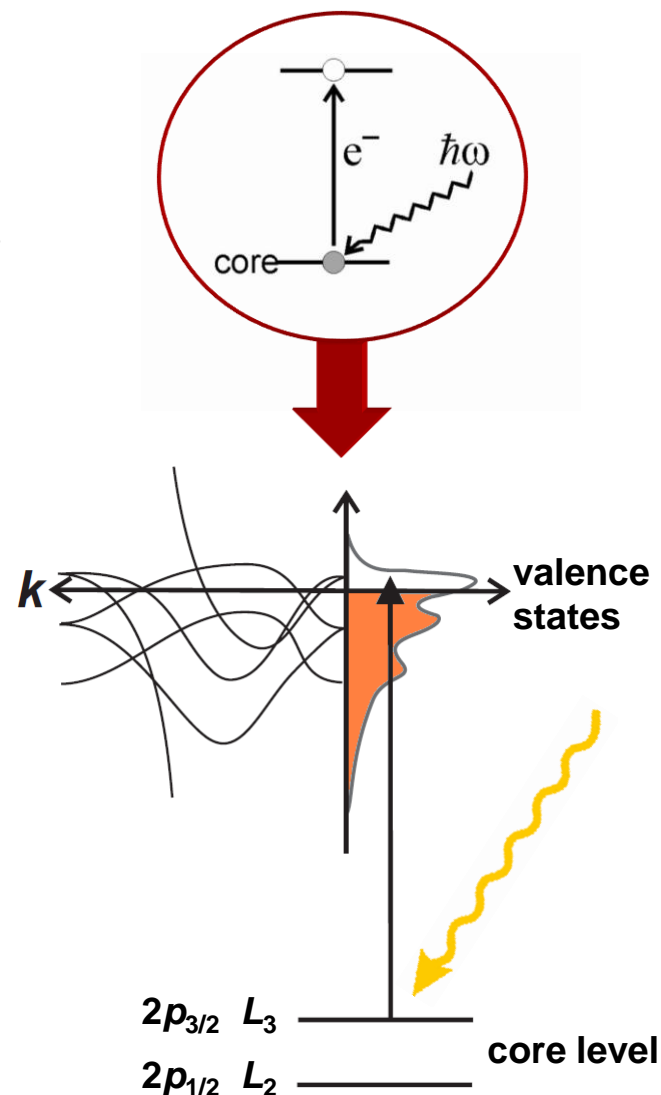
G. Van der Laan  
*Phys. Rev. B* **41**, 12366 (1990)

**X-ray absorption:**

- + Electrons excited from core shells to unoccupied valence states through absorption of a photon determined by energy and angular momentum conservation

**Simplest model: One electron picture**

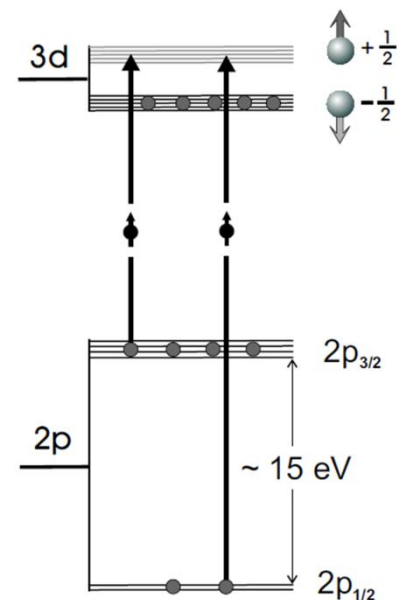
- + Photon transfers its energy and momentum to core electron
- + Core electron excited into unoccupied electronic state.
- + However: Not directly excited electrons also influenced by electron excitation, i.e. hole in core shell



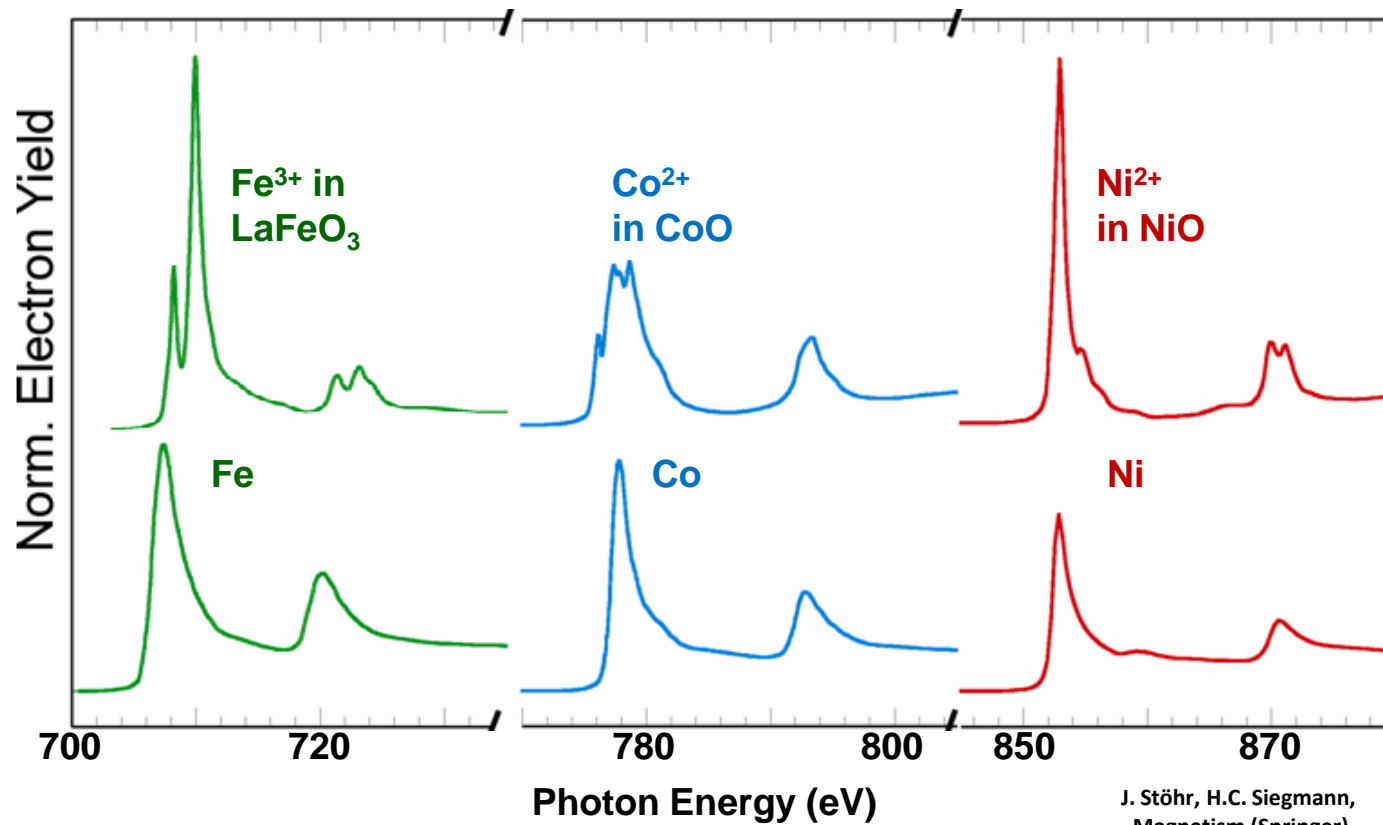


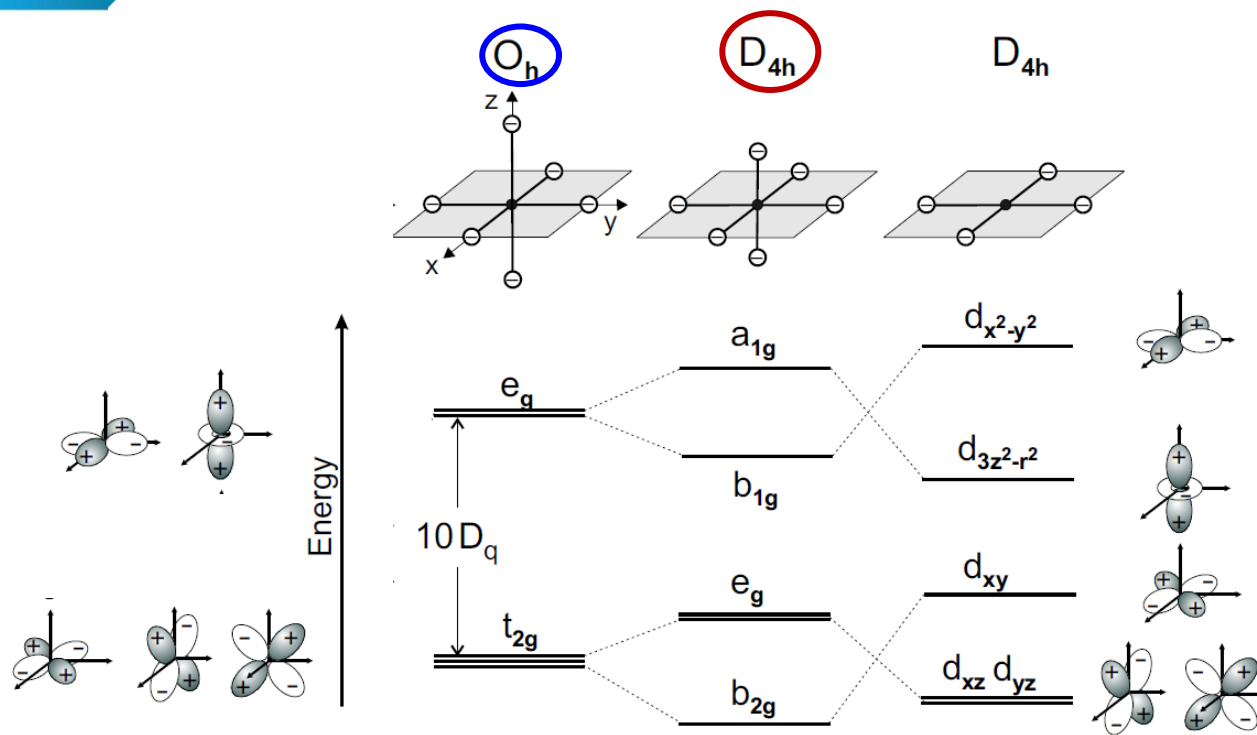
**Configuration model, e.g. *L* edge absorption :**

- + Atom is excited from ground/initial state configuration,  $2p^63d^n$  to excited/final state configuration,  $2p^53d^{n+1}$
- + Omission of all full subshells (spherical symmetric)
- + Takes into account correlation effects in the ground state as well as in the excited state
- + Leads to multiplet effects/structure

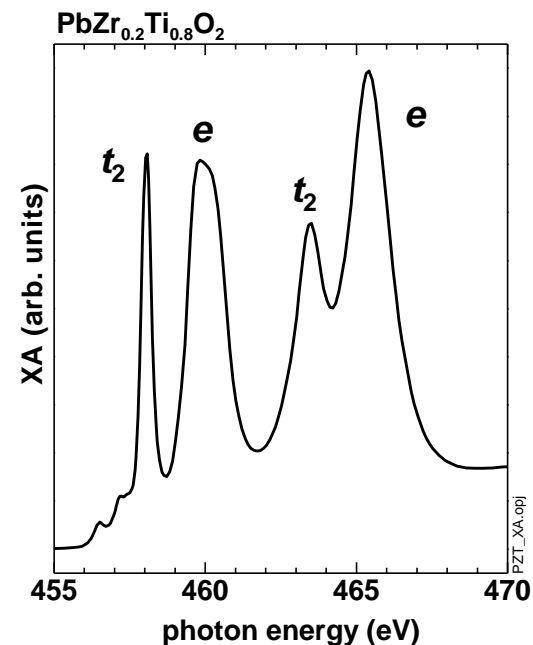


J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)





J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)



+ Electric dipole transitions:  $d^0 \rightarrow 2p^5 3d^1$

+ Crystal field splitting  $10Dq$  acting on 3d orbitals:

**Octahedral symmetry:**

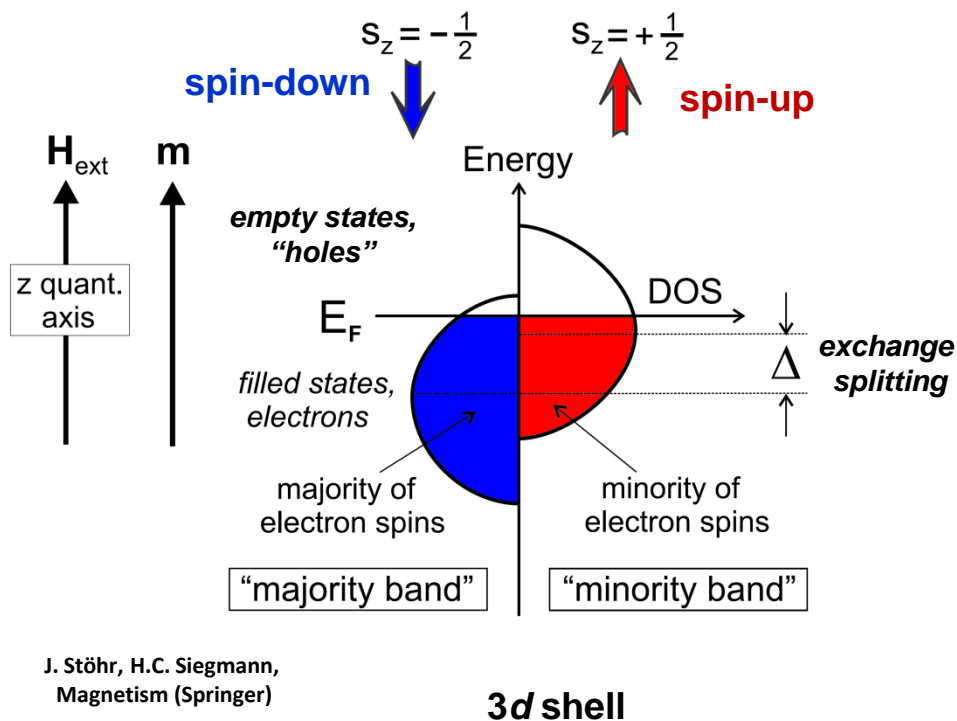
$e$  orbitals towards ligands  $\rightarrow$  higher energy

$t_2$  orbitals between ligands  $\rightarrow$  lower energy

**Tetragonal symmetry:**

$e$  orbitals  $\rightarrow b_2 = d_{xy}$ ,  $e = d_{yz}$ ,  $d_{yz}$

$t_2$  orbitals  $\rightarrow b_1 = d_{x^2-y^2}$ ,  $a_1 = d_{3z^2-r^2}$

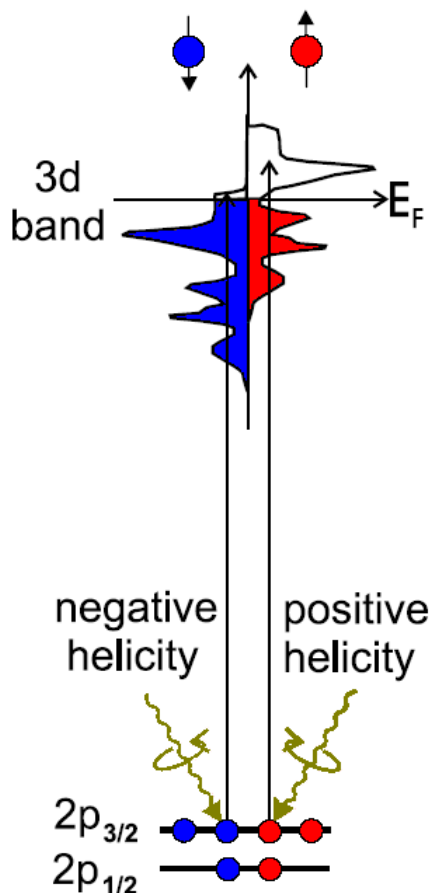


J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)

- + Magnetic moments in Fe, Co, Ni well described by Stoner model:  $d$ -bands containing up and down spins shifted relative to each other by "exchange splitting"
- + Spin- up and spin-down bands filled according to Fermi statistics
- + Magnetic moment  $|m|$  determined by difference in number of electrons in majority and minority bands

$$|m| = \mu_B (n_e^{maj} - n_e^{min})$$





J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)

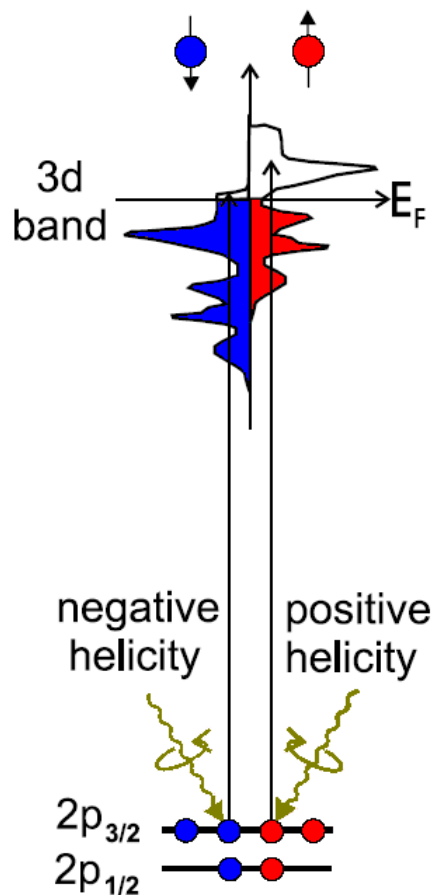
Photoelectrons excited from  $2p_{3/2}$ ,  $2p_{1/2}$  to  $3d$  states

### First step:

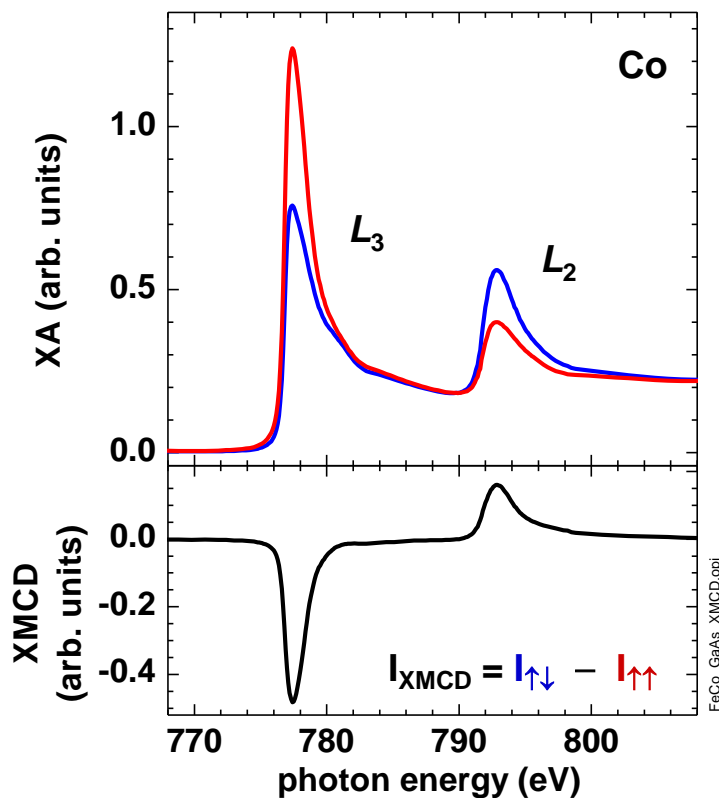
- + Excitation of electron from  $2p$  states by absorption of circularly polarized x rays
- + Note: Dipole operator does not affect spin  
 $\Rightarrow$  No spin flips during excitation
- + Conservation of angular momentum  
 $\Rightarrow$  Transfer of angular momentum ( $\pm\hbar$ ) from photon to electron
- + Spin-orbit coupling: Angular momentum of photon transferred in part to electron spin  
 $\Rightarrow$  Excited photoelectrons are spin polarized

### Second step:

- + Unequal spin-up and spin-down populations determine spin or orbital momentum of possible excitations



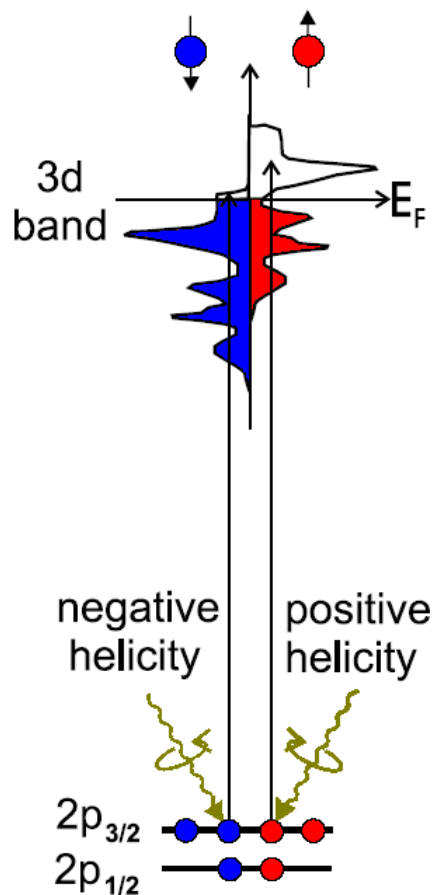
J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)



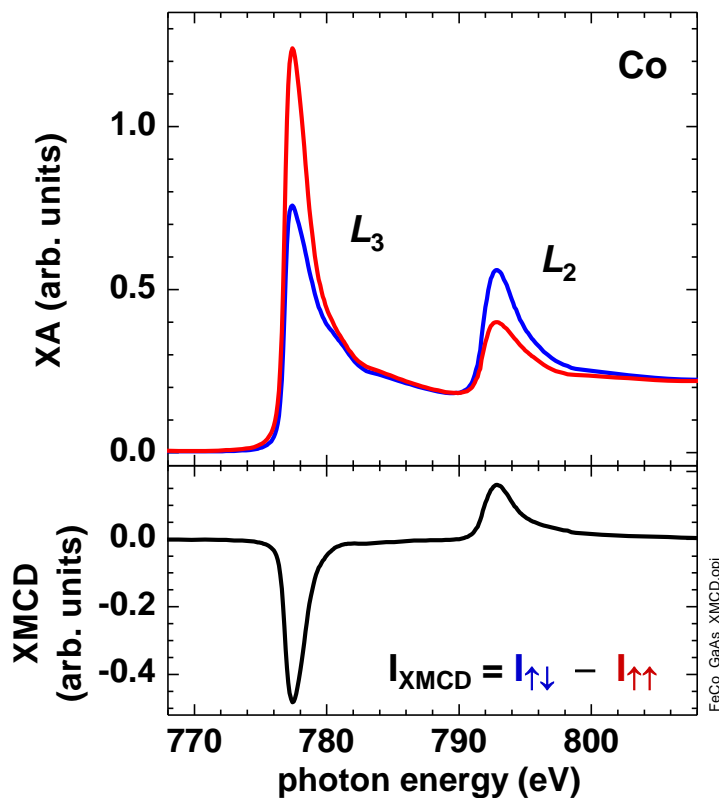
Magnitude of dichroism depends on

- + degree of circular photon polarization,  $P_{\text{circ}}$
- + angle  $\theta$  between photon angular momentum,  $L_{\text{ph}}$  and magnetic moment,  $m$
- + expectation value of 3d magnetic moment

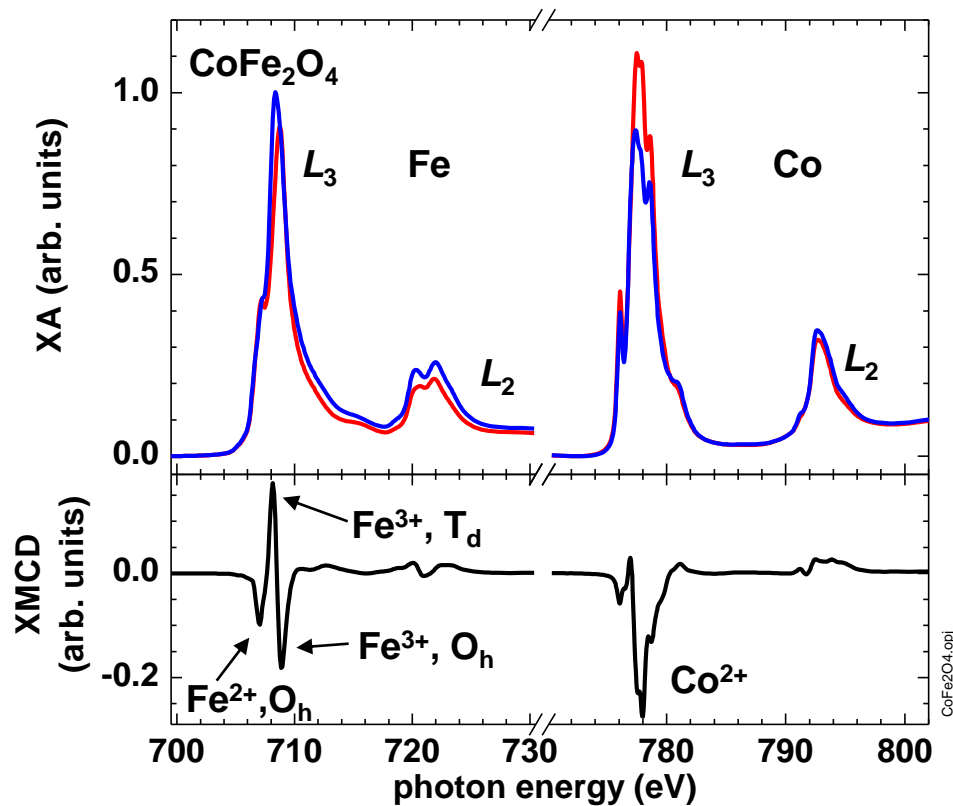
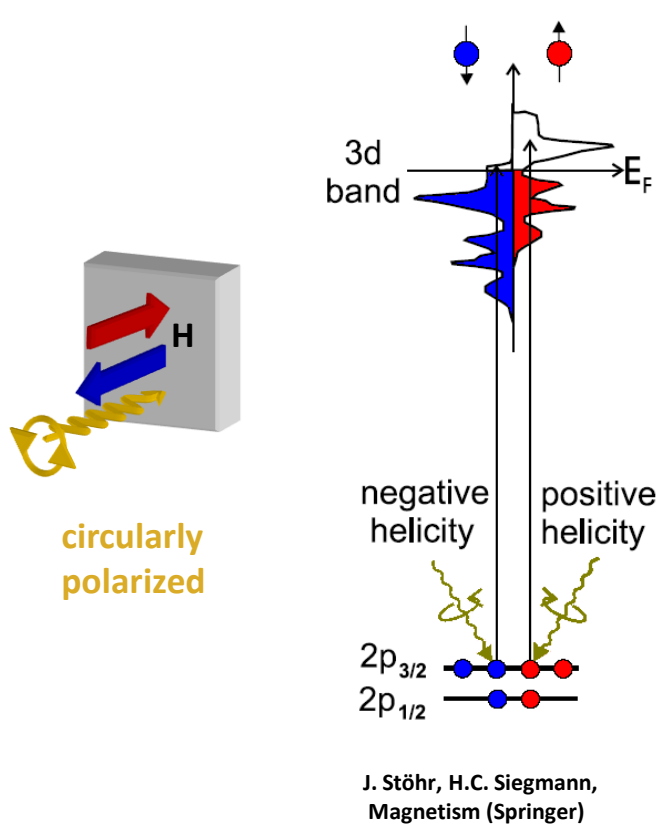
$$I_{XMCD} \propto P_{\text{circ}} \langle m \rangle \cos \theta$$



J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)



- +  $2p_{3/2}$  and  $2p_{1/2}$  have opposite spin orbit coupling ( $l+s$ ,  $l-s$ )  
 $\Rightarrow$  Spin polarization and XMCD have opposite sign at two edges
- + Spin polarization opposite for x rays with opposite helicity, i.e. photon spin,  $\pm\hbar$   
 $\Rightarrow$  XMCD reverses sign with polarization
- + Reversing x ray polarization is equivalent to reversing magnetization/spin direction



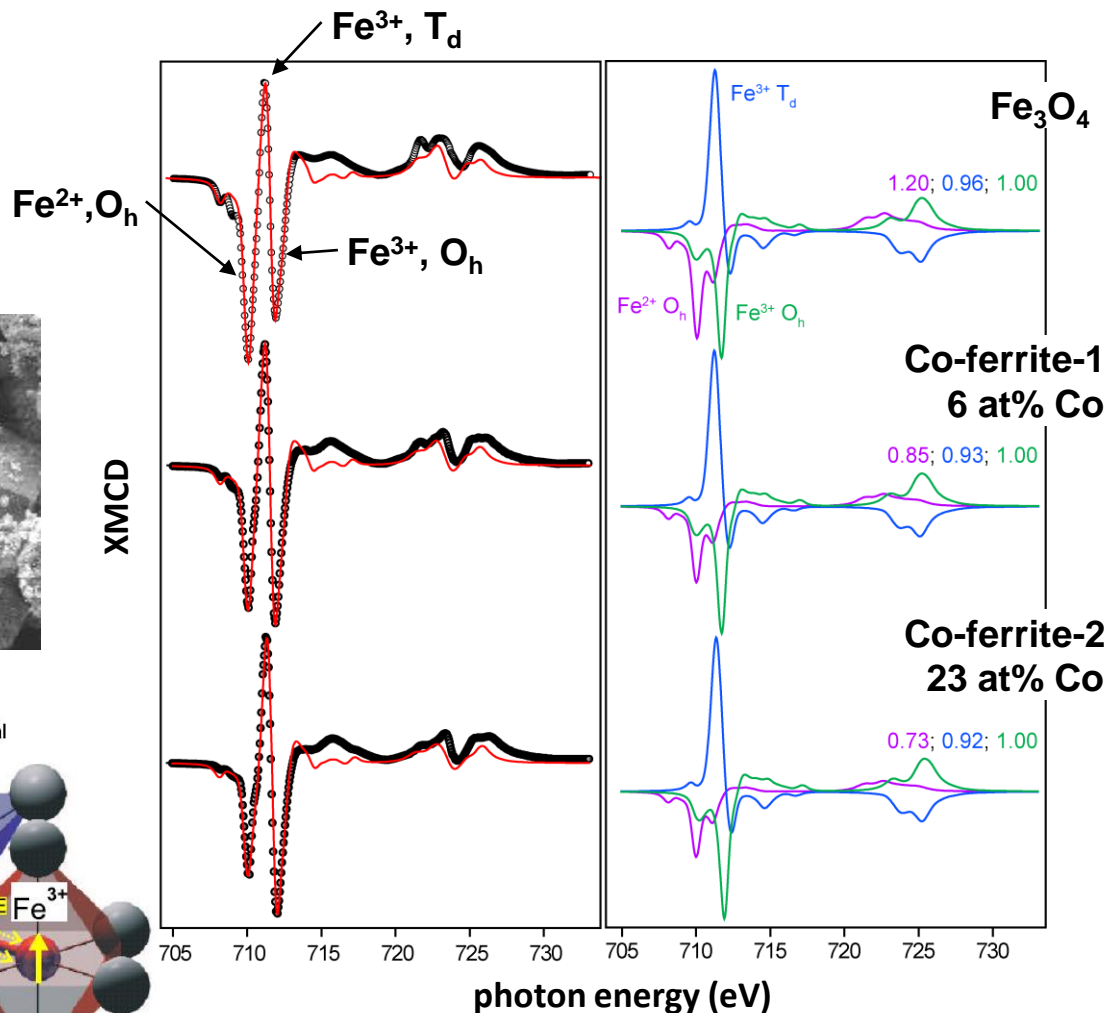
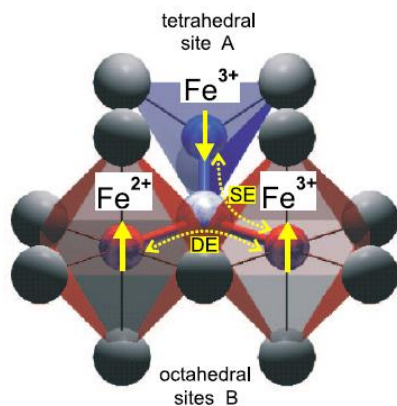
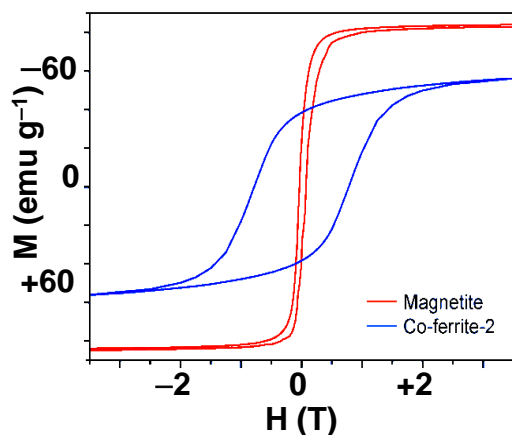
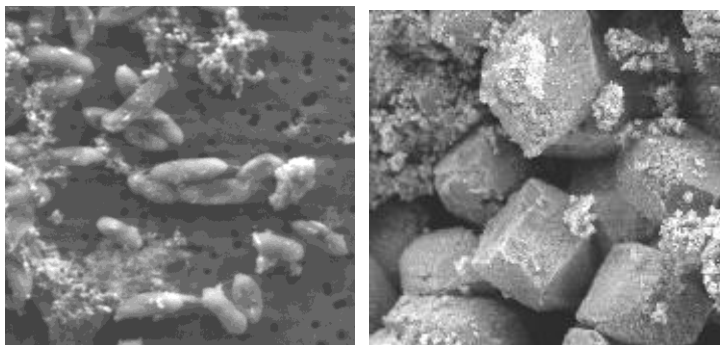
+ XMCD provides magnetic information resolving elements Fe, Co, ...

valence states: Fe<sup>2+</sup>, Fe<sup>3+</sup>, ...

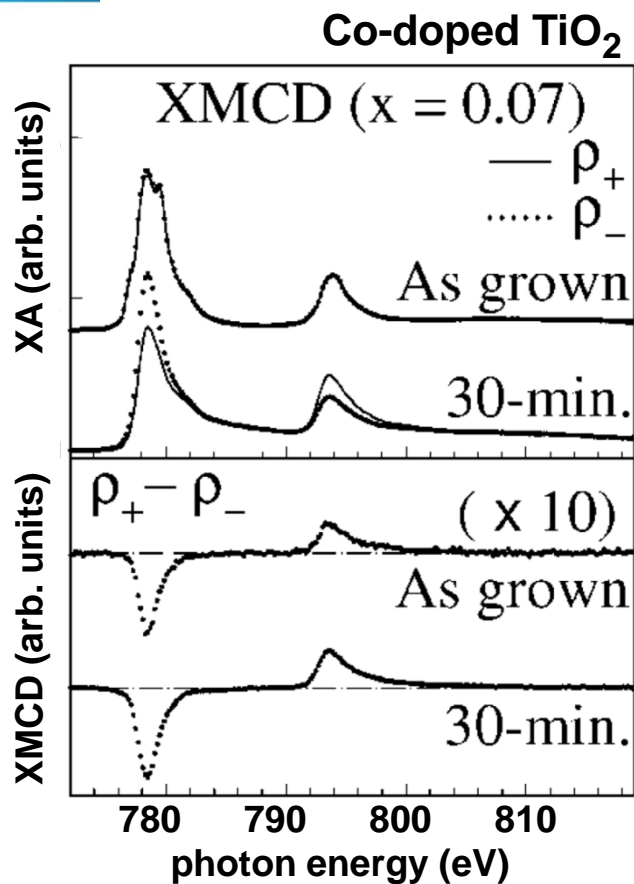
lattice sites: octahedral, O<sub>h</sub>, tetrahedral, T<sub>d</sub>, ...



+ *Geobacter sulfurreducens* bacteria form magnetite nanocrystals (15nm) via extracellular reduction of amorphous Fe(III)-bearing minerals

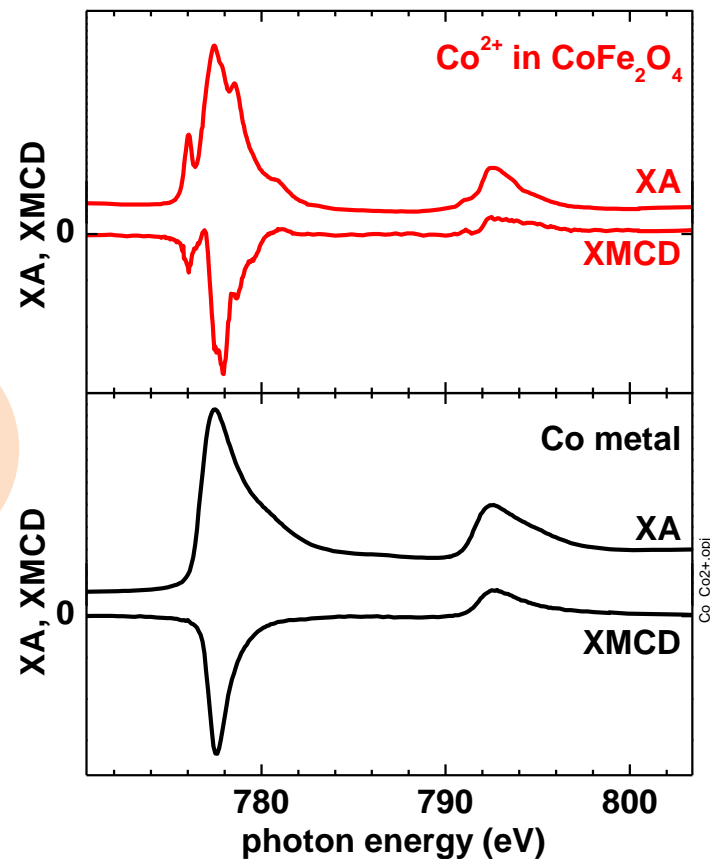
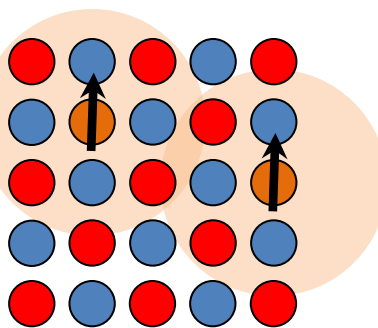


V. Cocker *et al.*,  
Eur. J. Mineral. **19**, 707–716 (2007)



J.-Y. Kim *et al.*,  
Phys. Rev. Lett. **90**, 017401 (2003)

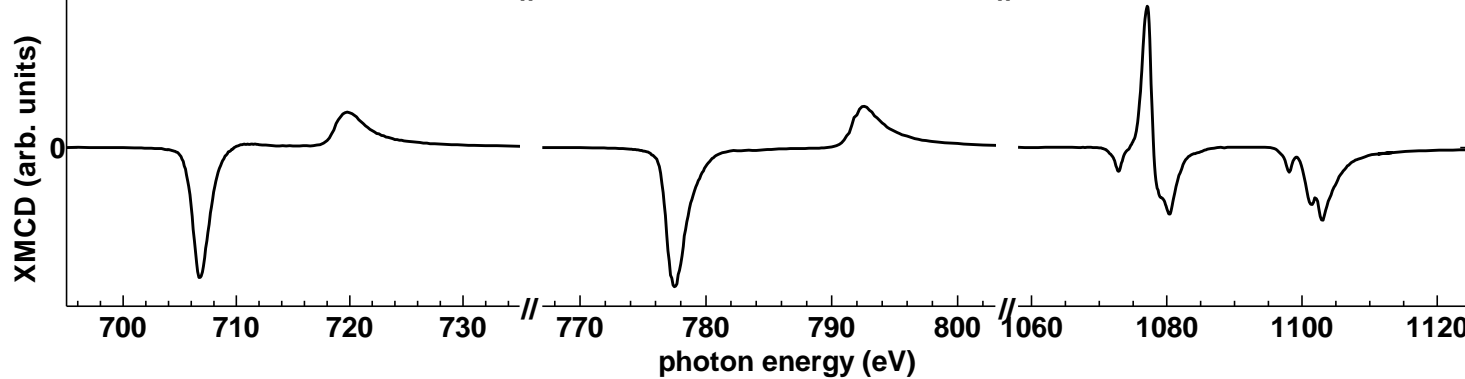
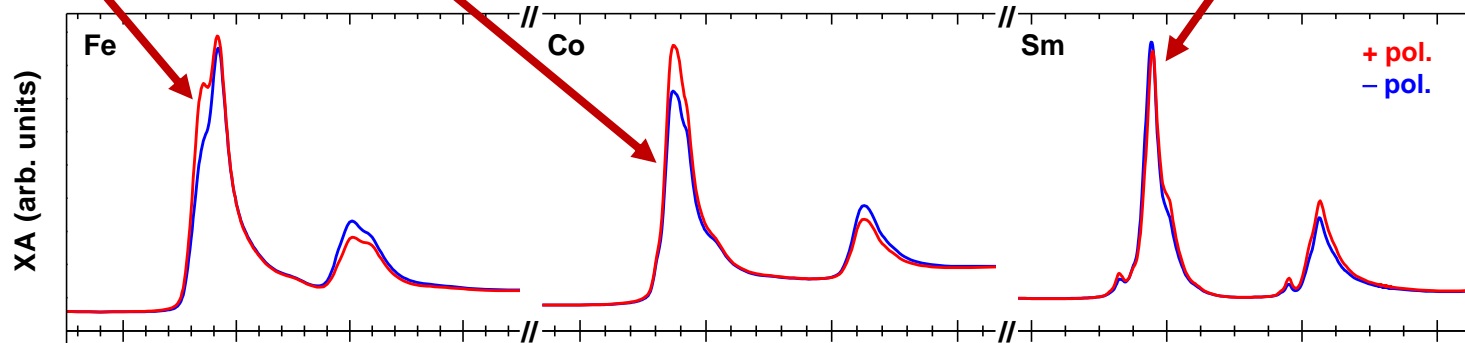
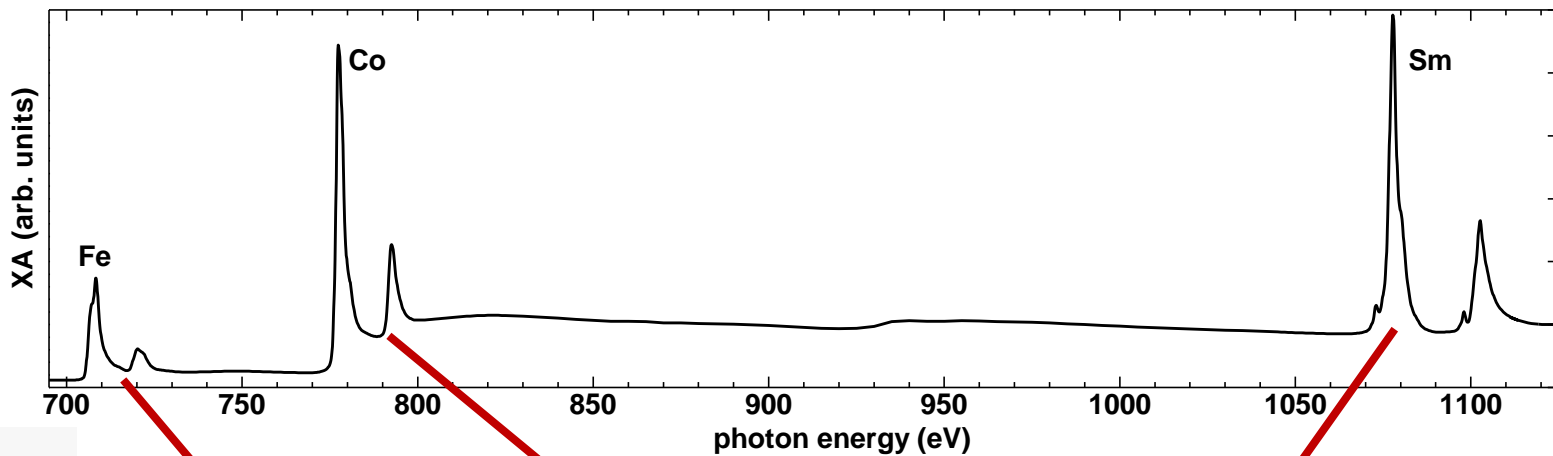
dilute magnetic  
semiconductors

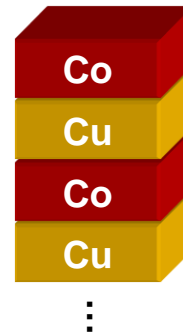
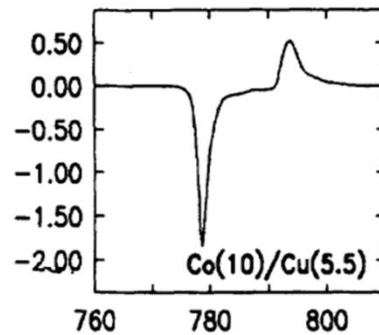
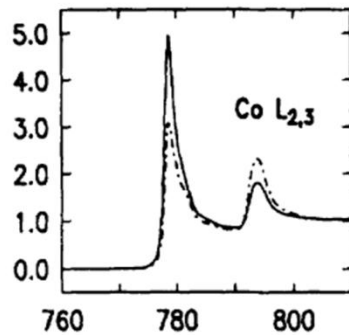


+ Comparing XMCD spectra with model compounds and/or calculations

⇒ Identifying magnetic phases

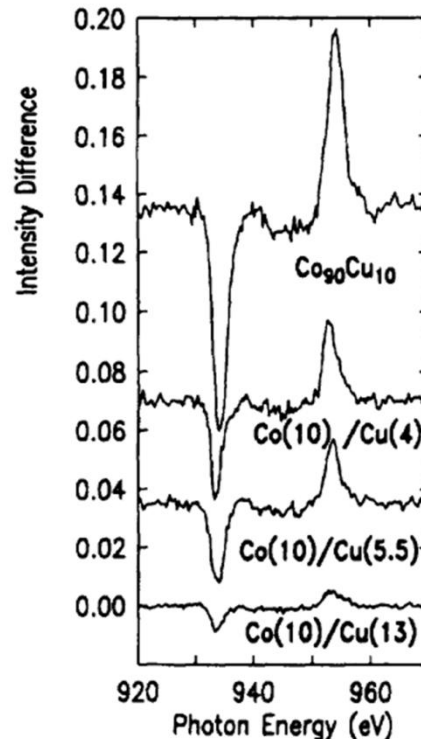
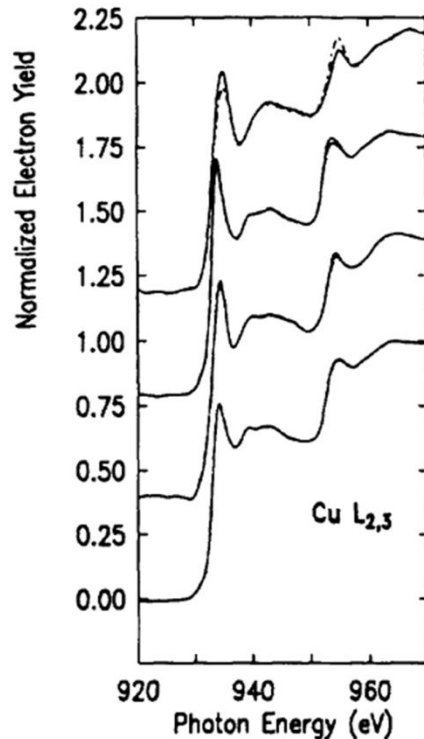
'CoSm'



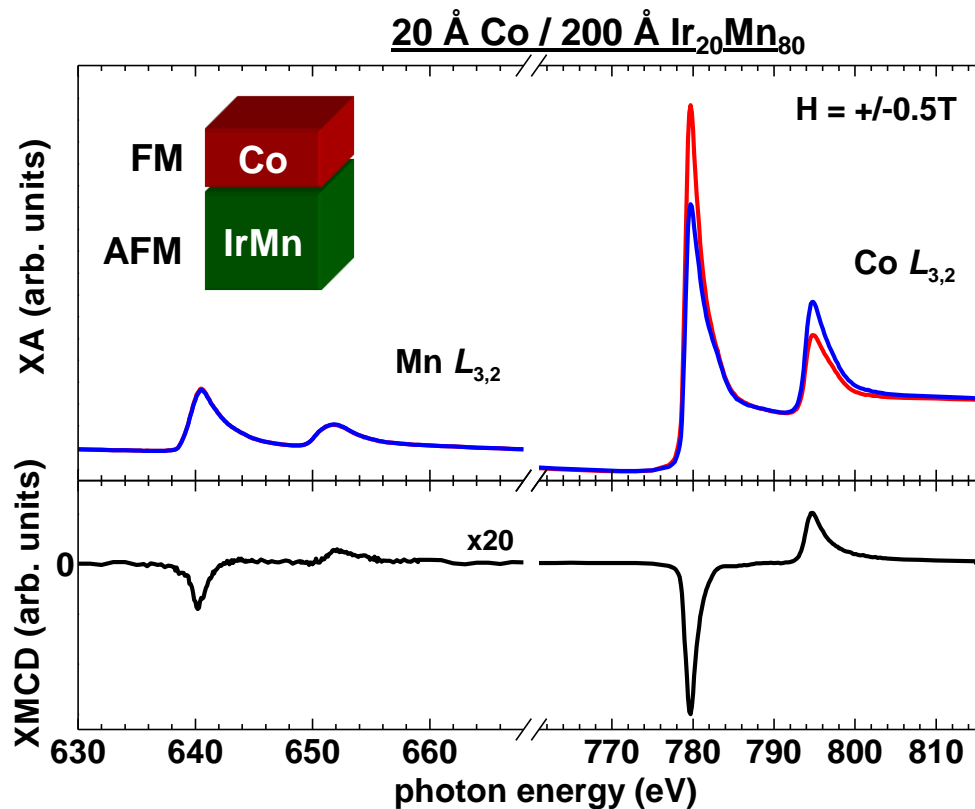


- + The element-specificity makes XMCD measurements an ideal tool to determine induced moments at interfaces between magnetic and non-magnetic elements.

M. G. Samant *et al.*,  
Phys. Rev. Lett. 72, 1112 (1994)

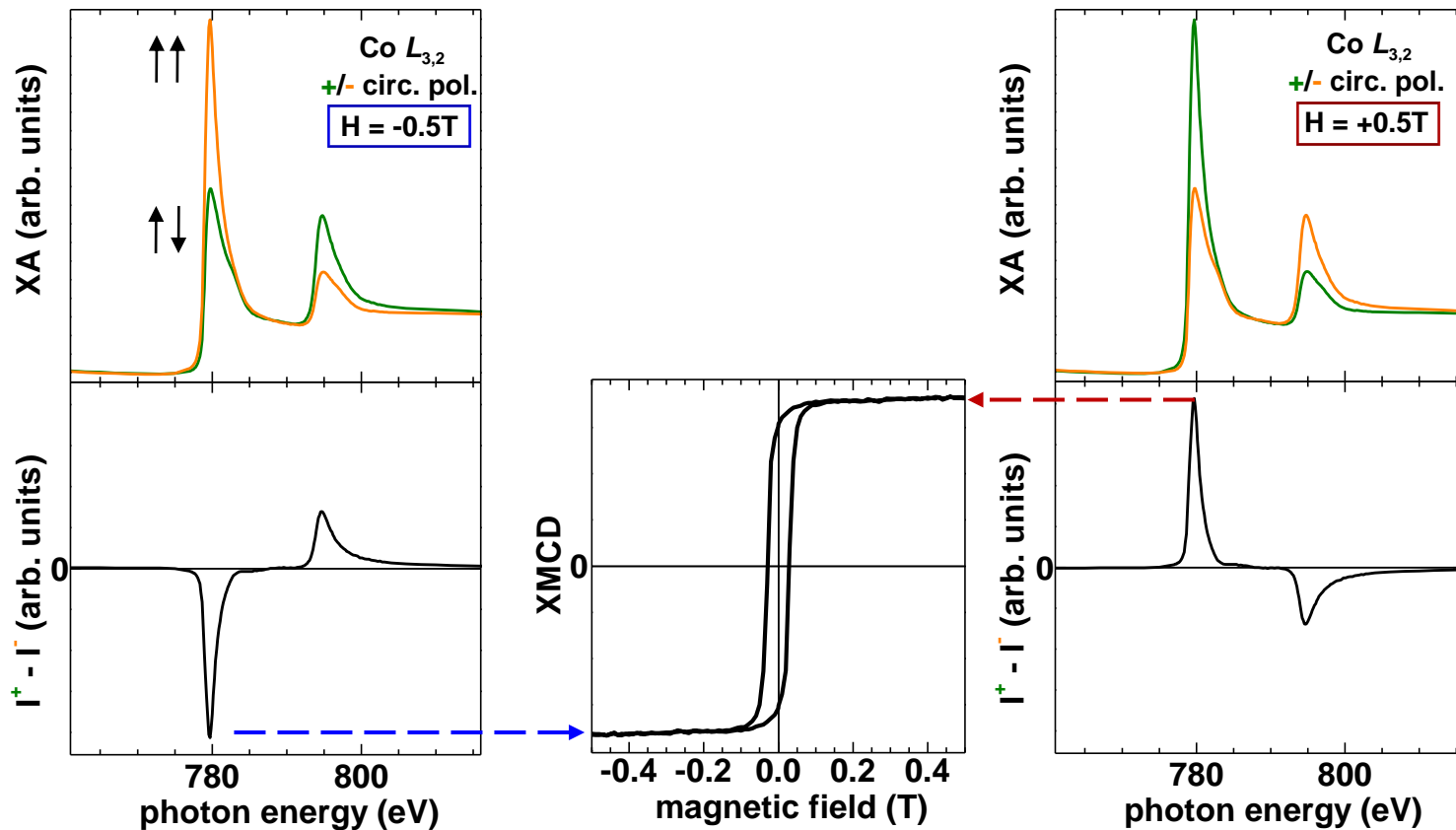


- + Weak Mn XMCD signal
- ⇒ Uncompensated Mn at Co/IrMn interface
- + Same sign of XMCD signal for Co and Mn
- ⇒ Parallel coupling of Co and Mn moments
- + Nominal thickness of uncompensated interface moments:  $(0.5 \pm 0.1)$  ML for Co/Ir<sub>20</sub>Mn<sub>80</sub>.

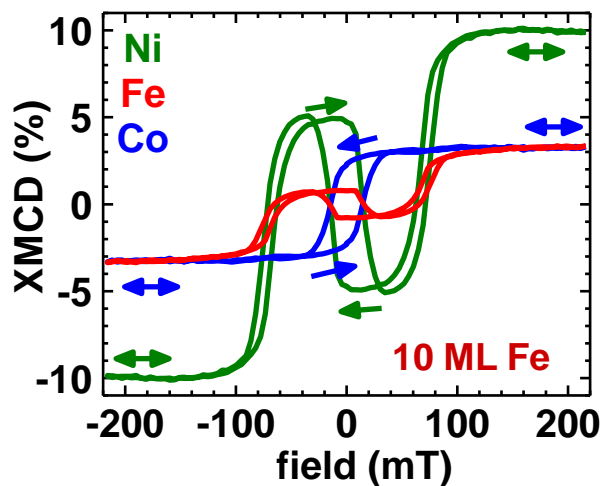
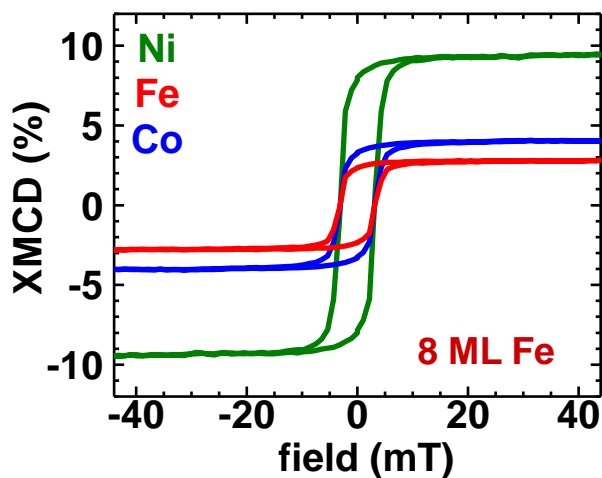
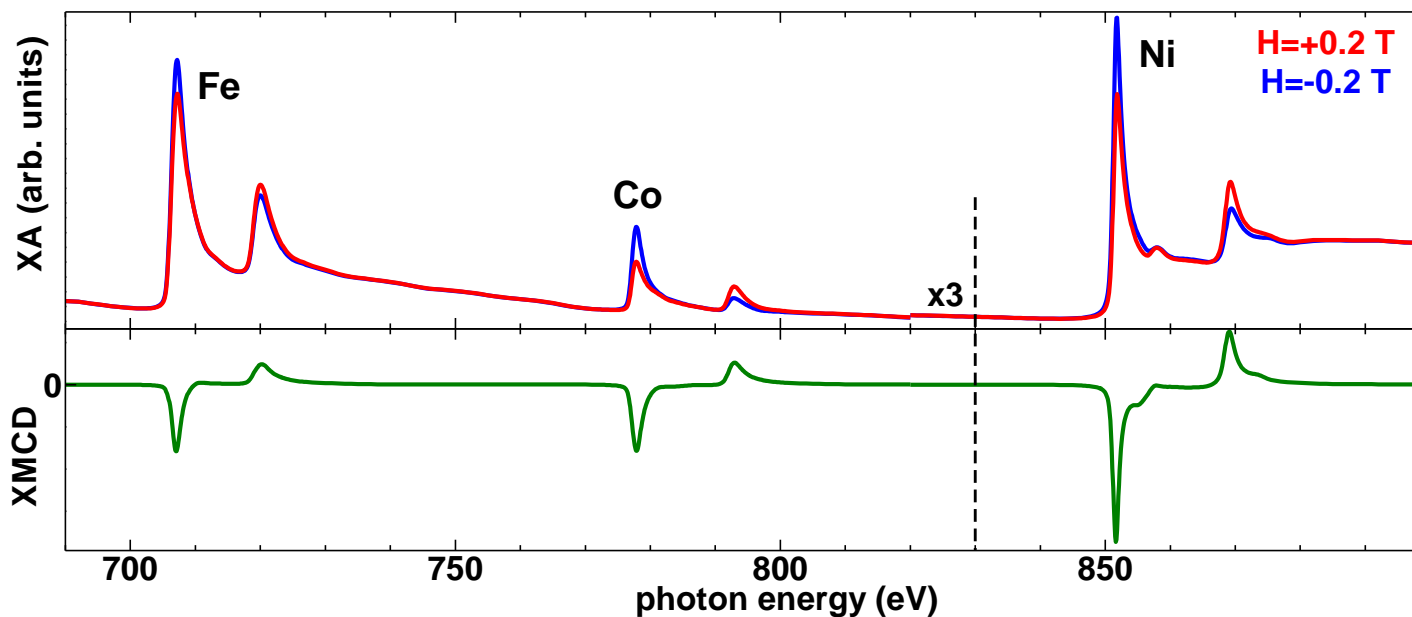
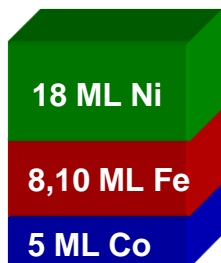


H. Ohldag *et al.*,  
Phys. Rev. Lett. 91, 017203 (2003)

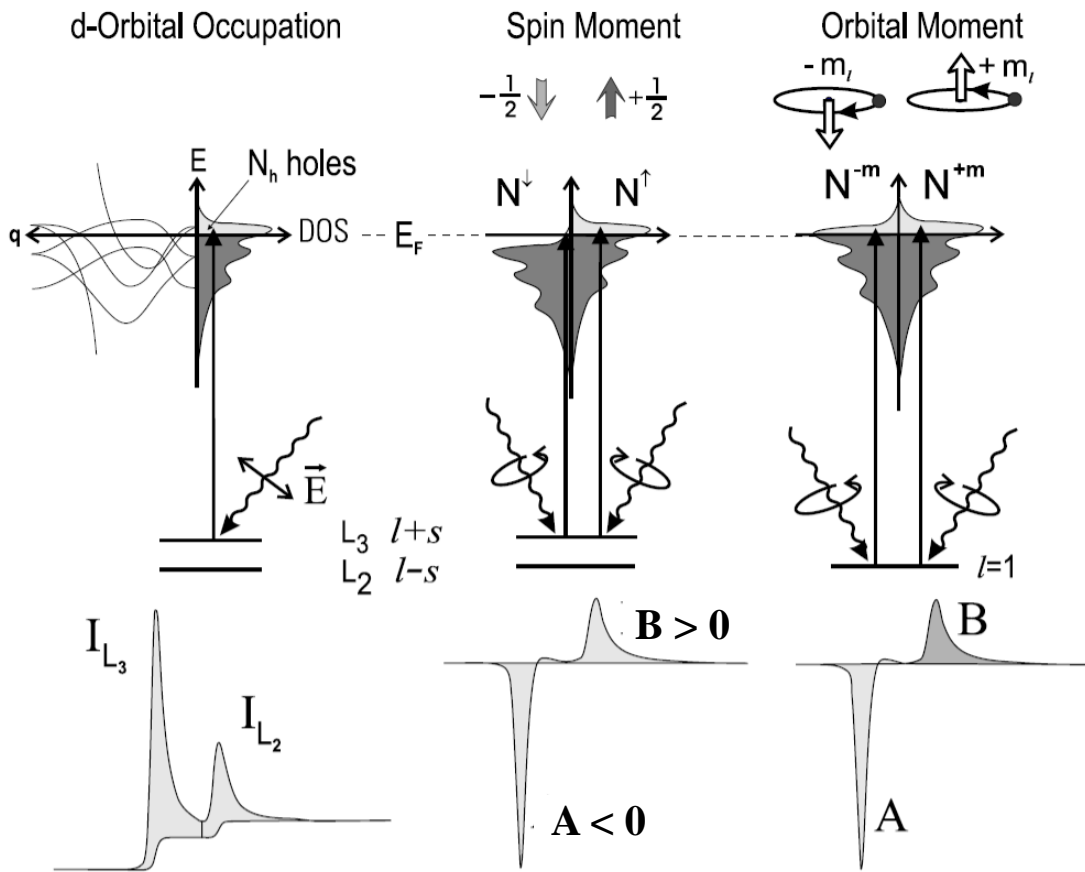




- + Monitoring field dependence of XMCD
- ⇒ Element-specific information on magnetization reversal in complex magnetic nanostructures.



- + Monitoring field dependence of XMCD
- ⇒ Detailed information on magnetization reversal in complex magnetic heterostructures



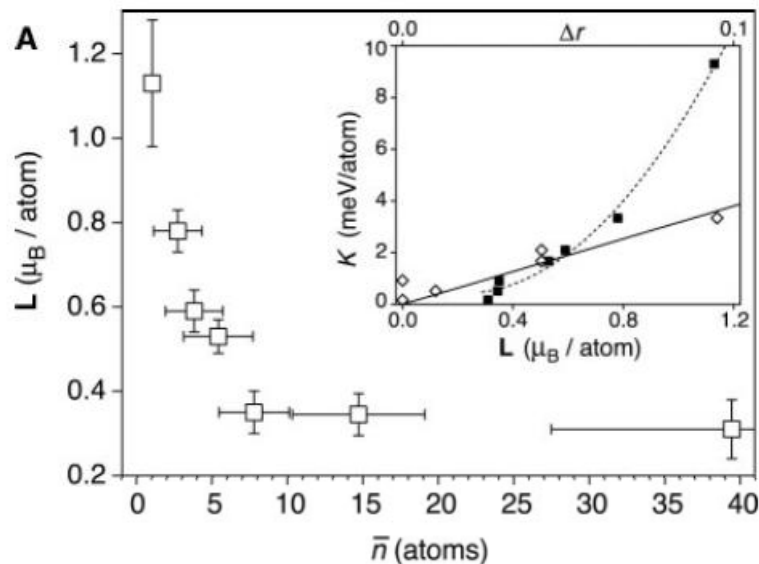
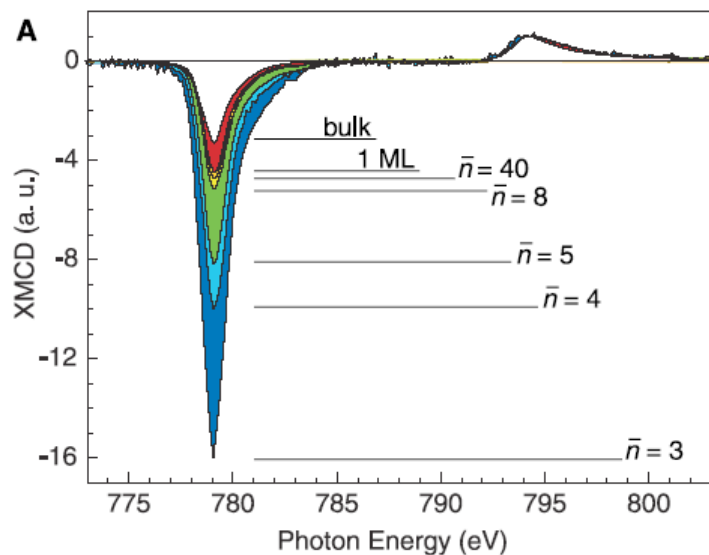
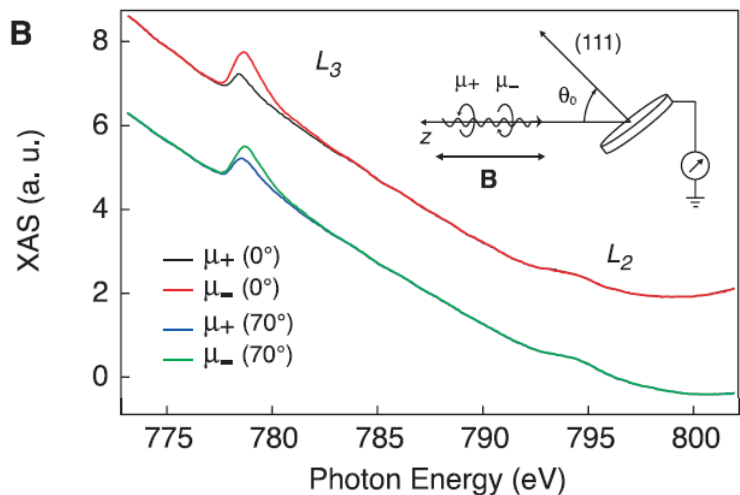
**+ Theoretically derived sum rules correlate XMCD spectra with spin and orbital moment providing unique tool for studying magnetic materials.**

$$N_h = \langle I_{L3} + I_{L2} \rangle / C$$

$$m_L = -2\mu_B \langle A + B \rangle / 3C$$

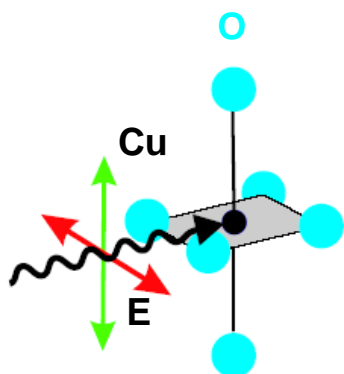
$$m_S = \mu_B \langle -A + 2B \rangle / C$$

J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)

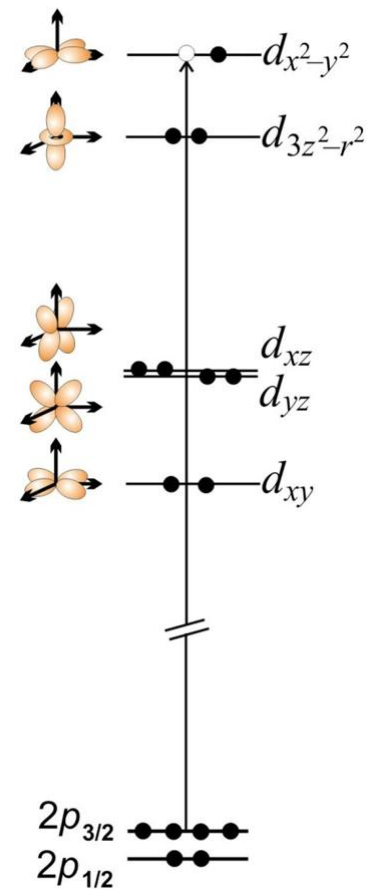
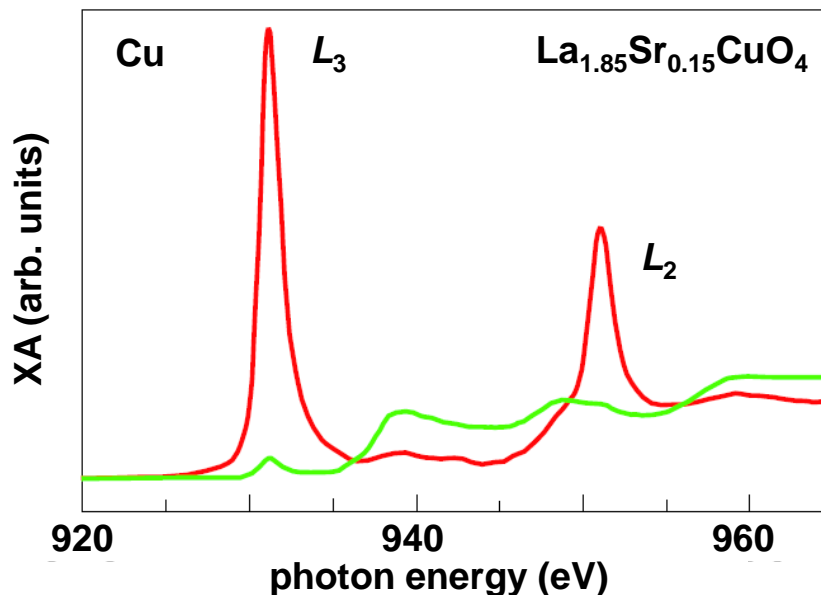


- + Strong variation of orbital and spin magnetic moment observable as change in relative  $L_3$  and  $L_2$  intensity in XMCD spectrum.
- + Co atoms and nanoparticles on Pt have enhanced orbital moments up to  $1.1 \mu_B$

P. Gambardella *et al.*,  
 Science **300**, 1130 (2003)



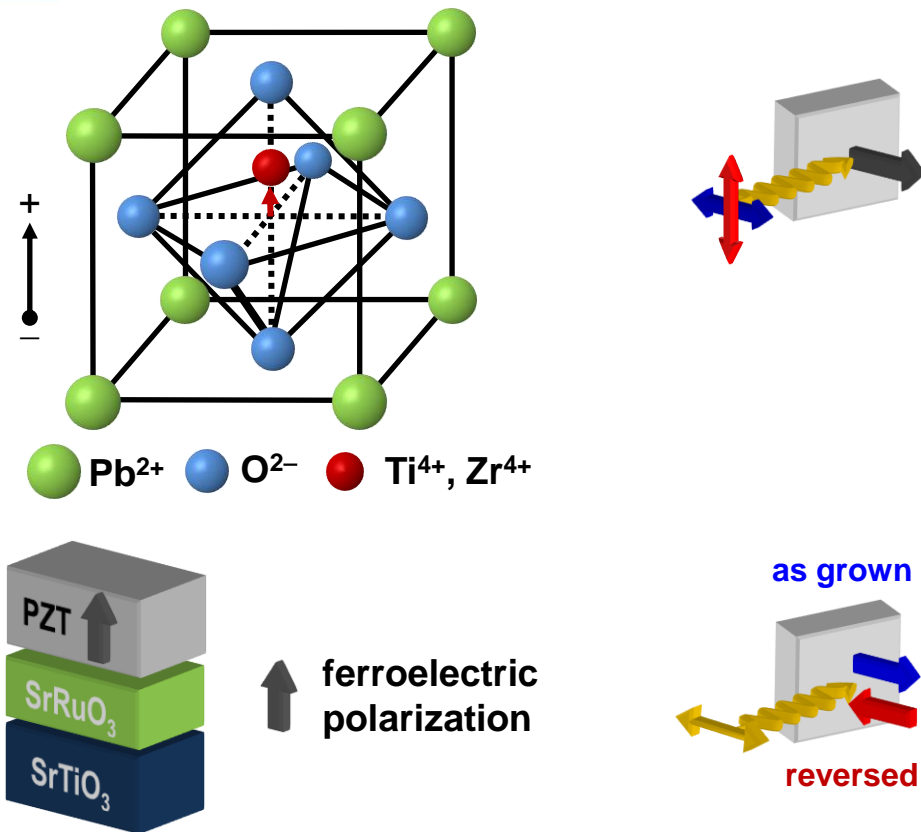
C. T. Chen *et al.*  
PRL **68**, 2543 (1992)



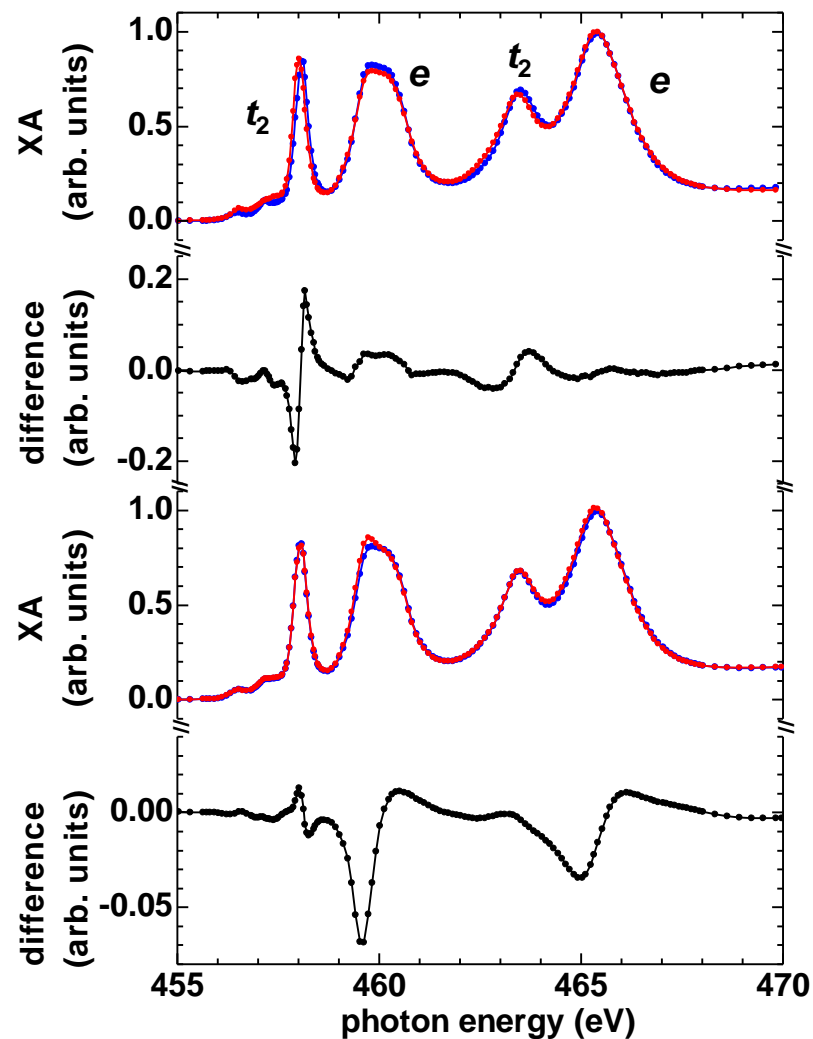
### X-Ray Linear Dichroism:

- + Difference in x-ray absorption for different linear polarization direction relative to crystalline and/or spin axis.
- + Due to the anisotropic charge distribution about the absorbing atom caused by bonding and/or magnetic order.
- + “Search Light Effect”: X-ray absorption of linear polarized x rays proportional to density of empty valence states in direction of electric field vector E.

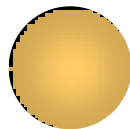




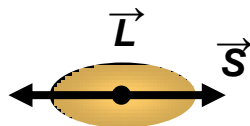
- + Spontaneous electric polarization due to off-center shift of  $\text{Ti}^{4+}, \text{Zr}^{4+}$  associated with tetragonal distortion  $\Leftrightarrow$  linear dichroism
- + Reversing ferroelectric polarization changes XA  $\Leftrightarrow$  Change in tetragonal distortion



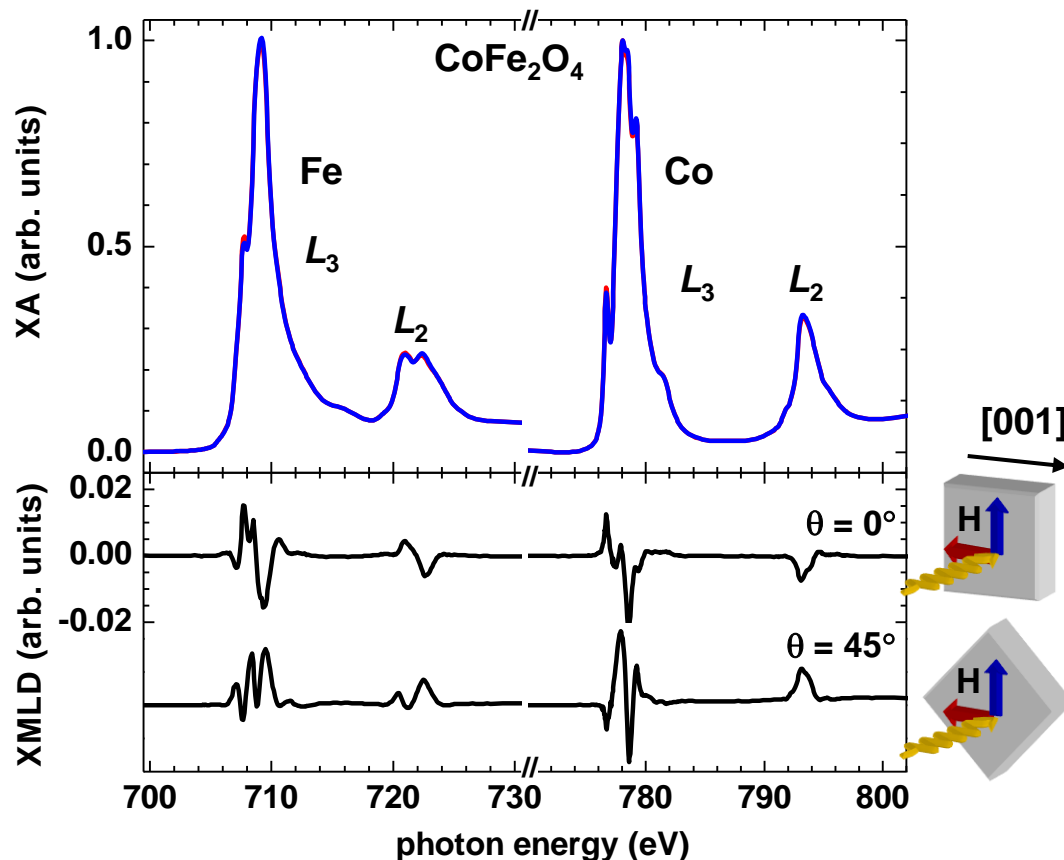
E. Arenholz *et al.*,  
 Phys. Rev. B **82**, 140103 (2010)



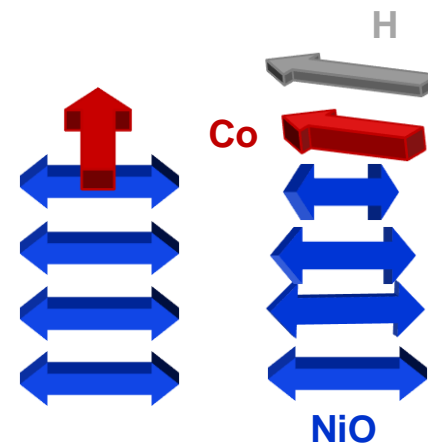
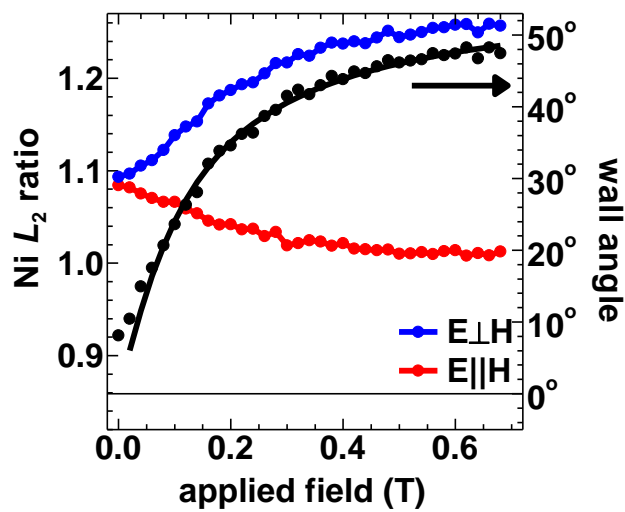
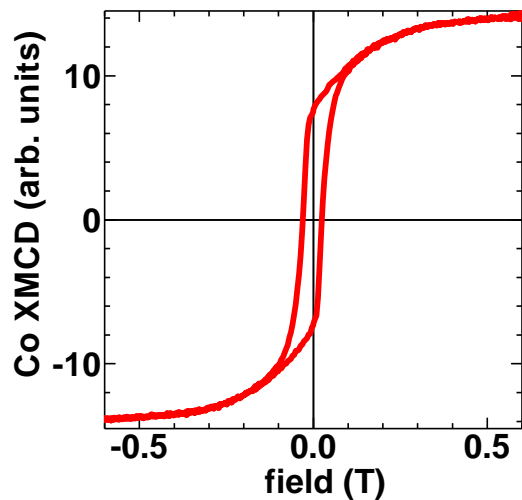
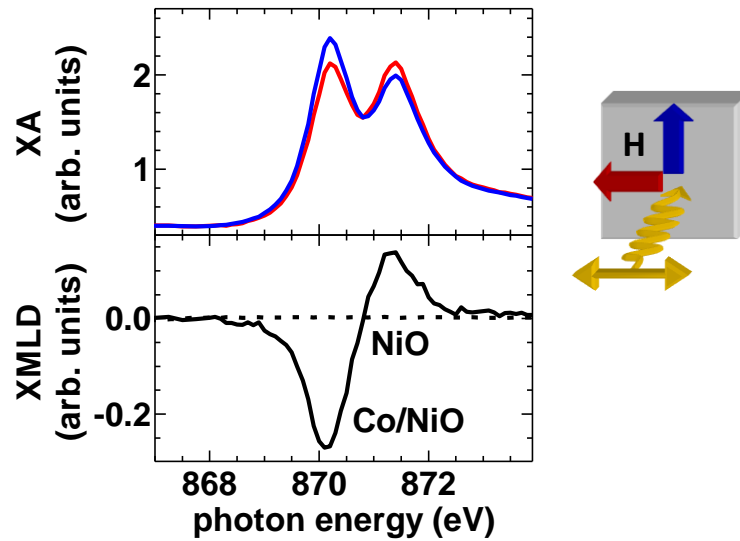
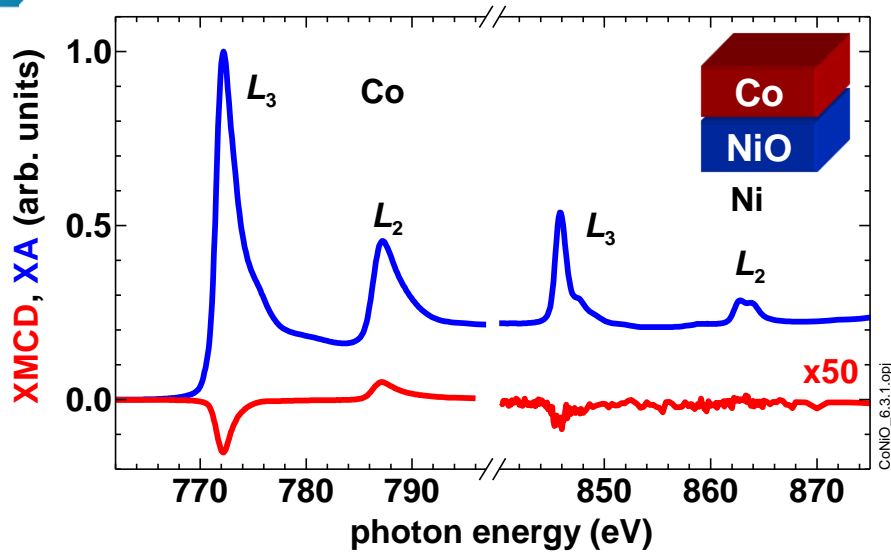
Isotropic  $d$  electron charge density  
 $\Rightarrow$  No polarization dependence



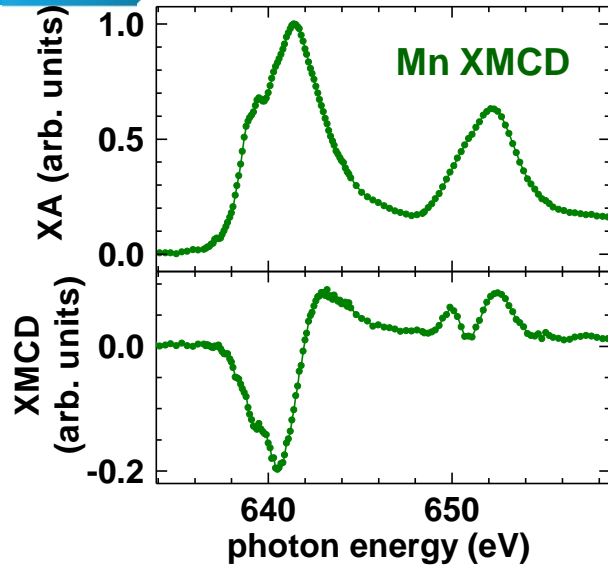
Magnetically aligned system  
 $\Rightarrow$  Spin-orbit coupling distorts  
 charge density  
 $\Rightarrow$  Polarization dependence



- +  $I_{\text{XMLD}} = I_{\parallel} - I_{\perp} \propto \langle m^2 \rangle$ ,  $\langle m^2 \rangle =$  expectation value of the square of the atomic magnetic moment
- + XMLD allows investigating ferri- and ferromagnets as well as antiferromagnets
- + XMLD spectral shape and angular dependence are determined by magnetic order and lattice symmetry

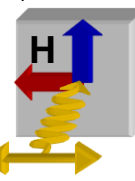
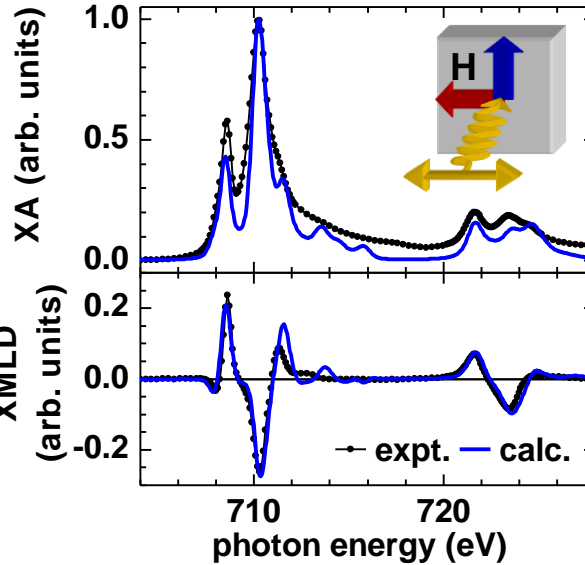


A. Scholl *et al.*,  
Phys. Rev. Lett. **92**, 247201 (2004)

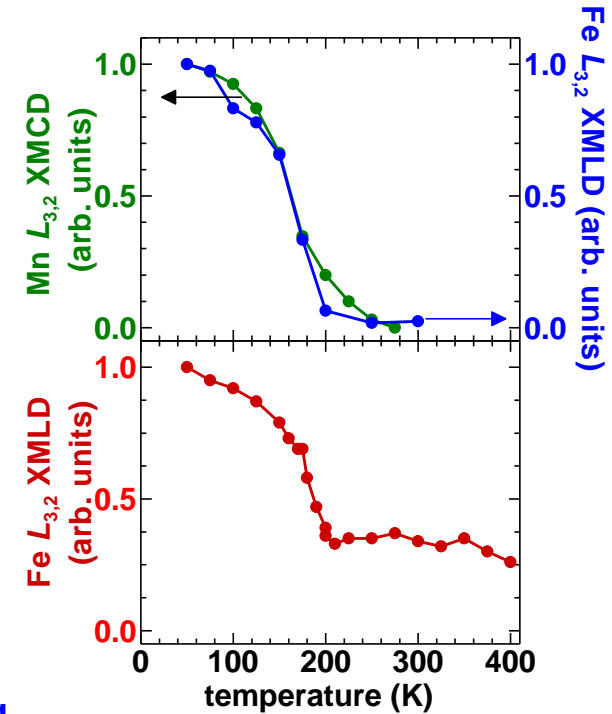
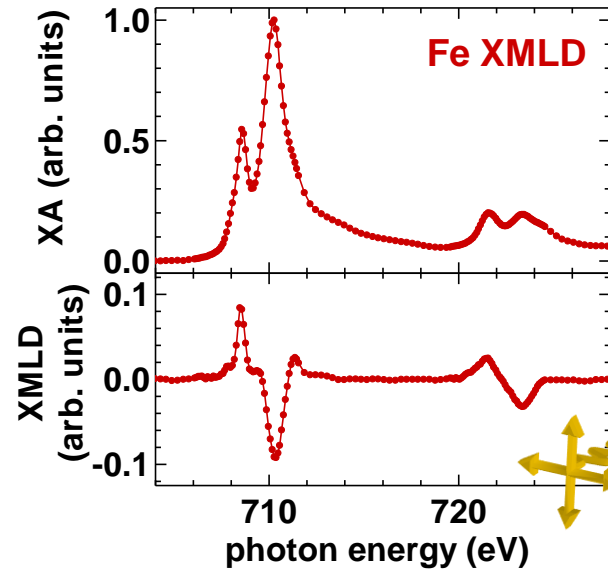


$\text{La}_{0.7}\text{Sr}_{0.3}\text{MnO}_3$  (LSMO)  
ferromagnet

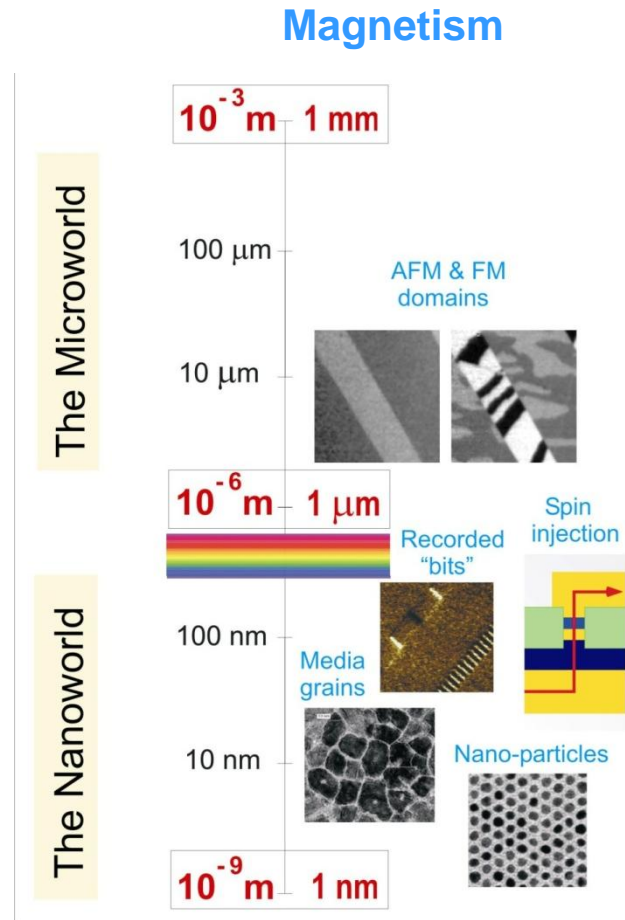
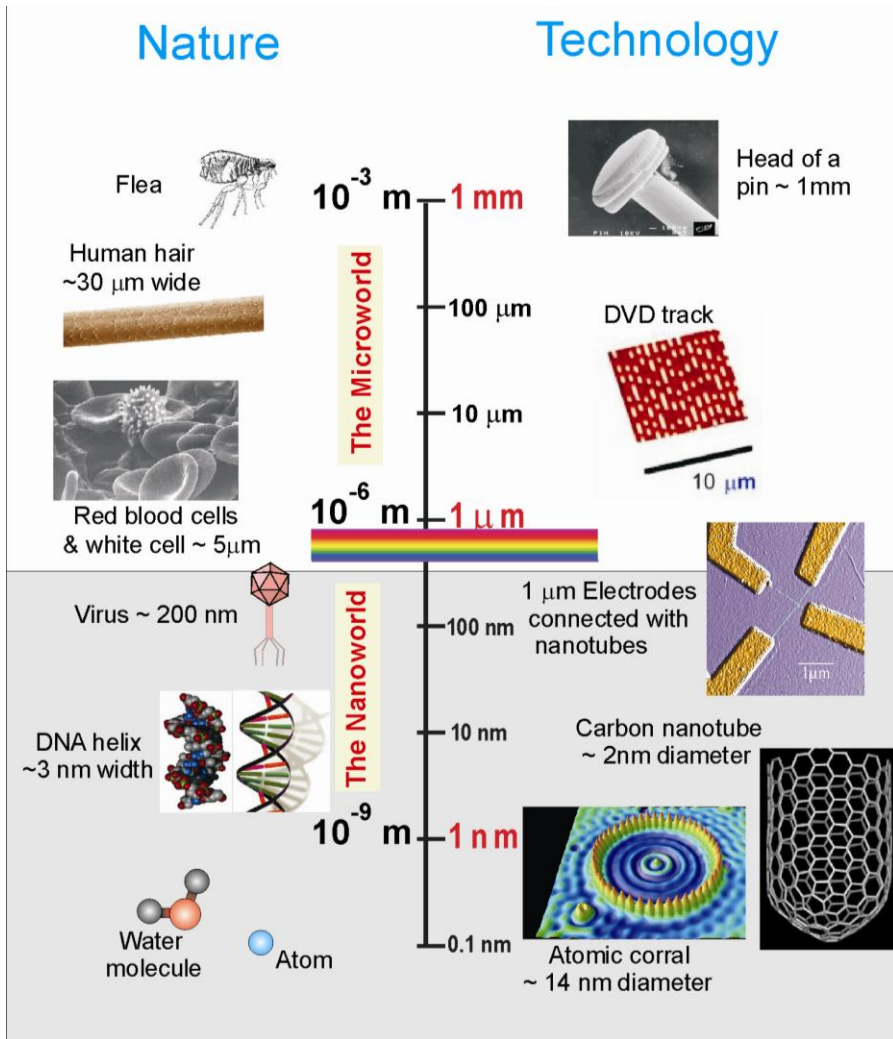
$\text{La}_{0.7}\text{Sr}_{0.3}\text{FeO}_3$  (LSFO)  
antiferromagnet



⇒ Perpendicular coupling  
at LSMO/LSFO interface

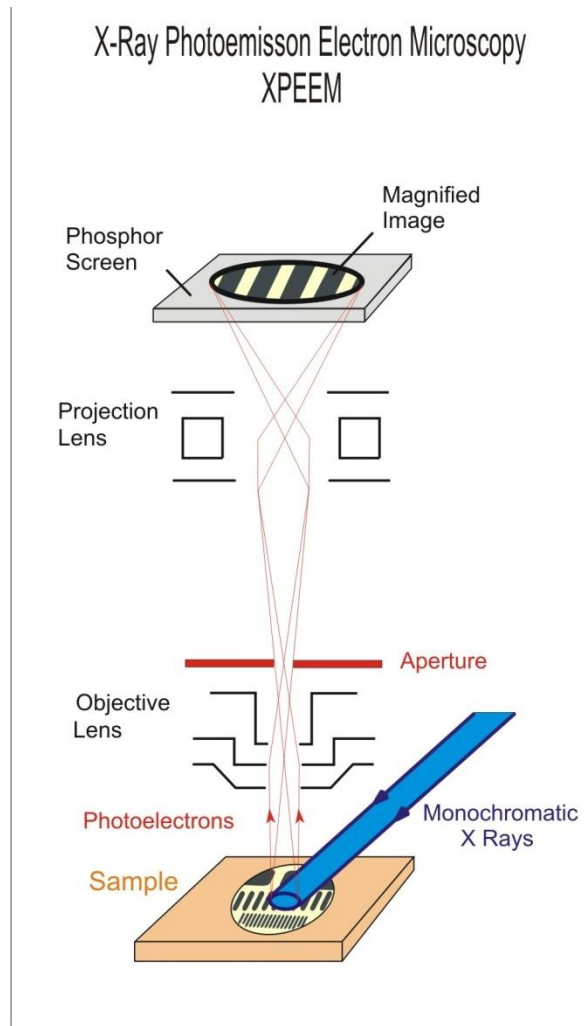
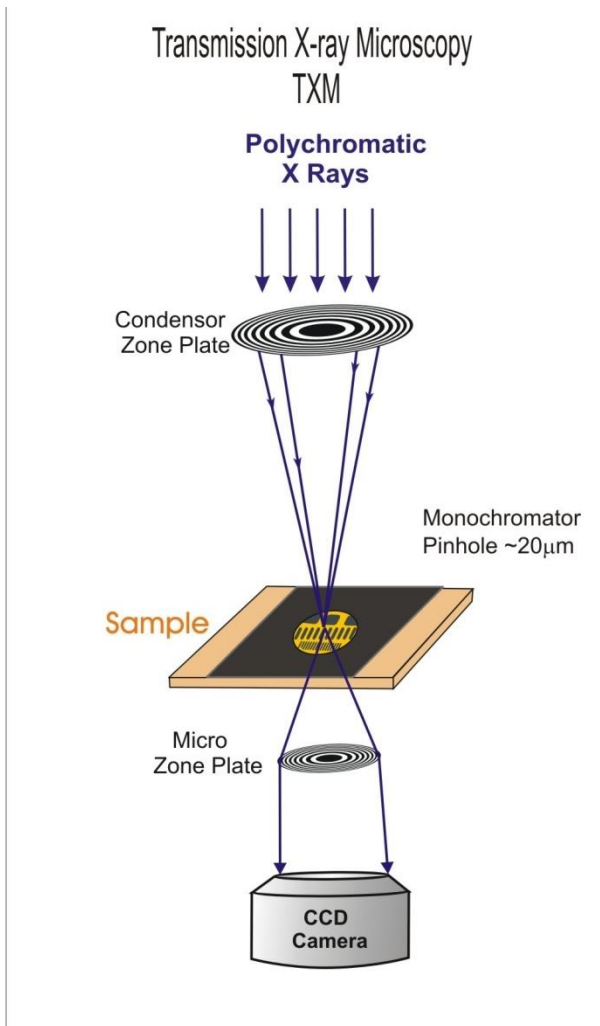
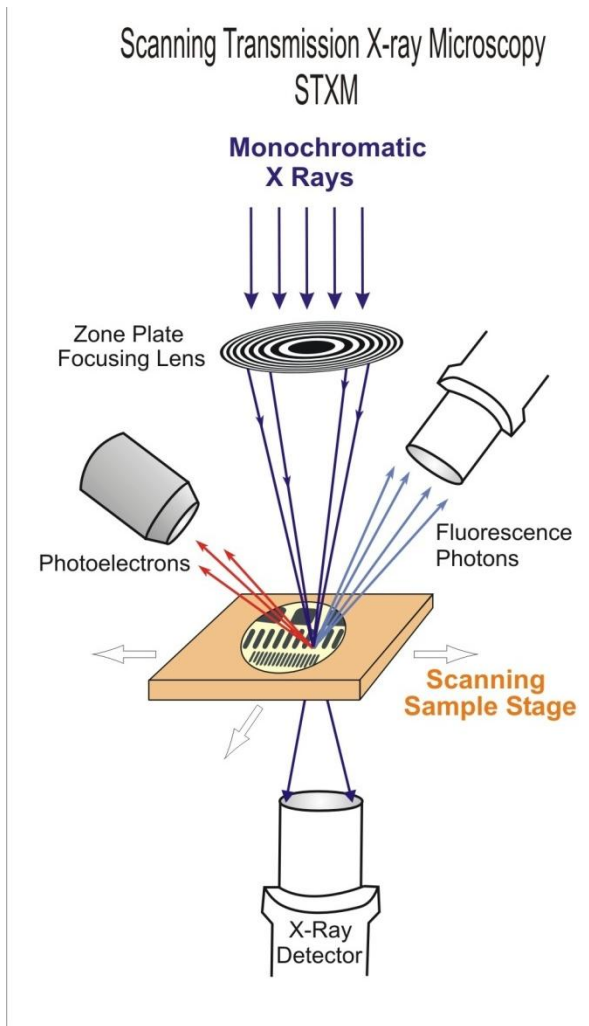


E. Arenholz *et al.*,  
Appl. Phys. Lett. **94**, 072503 (2009)



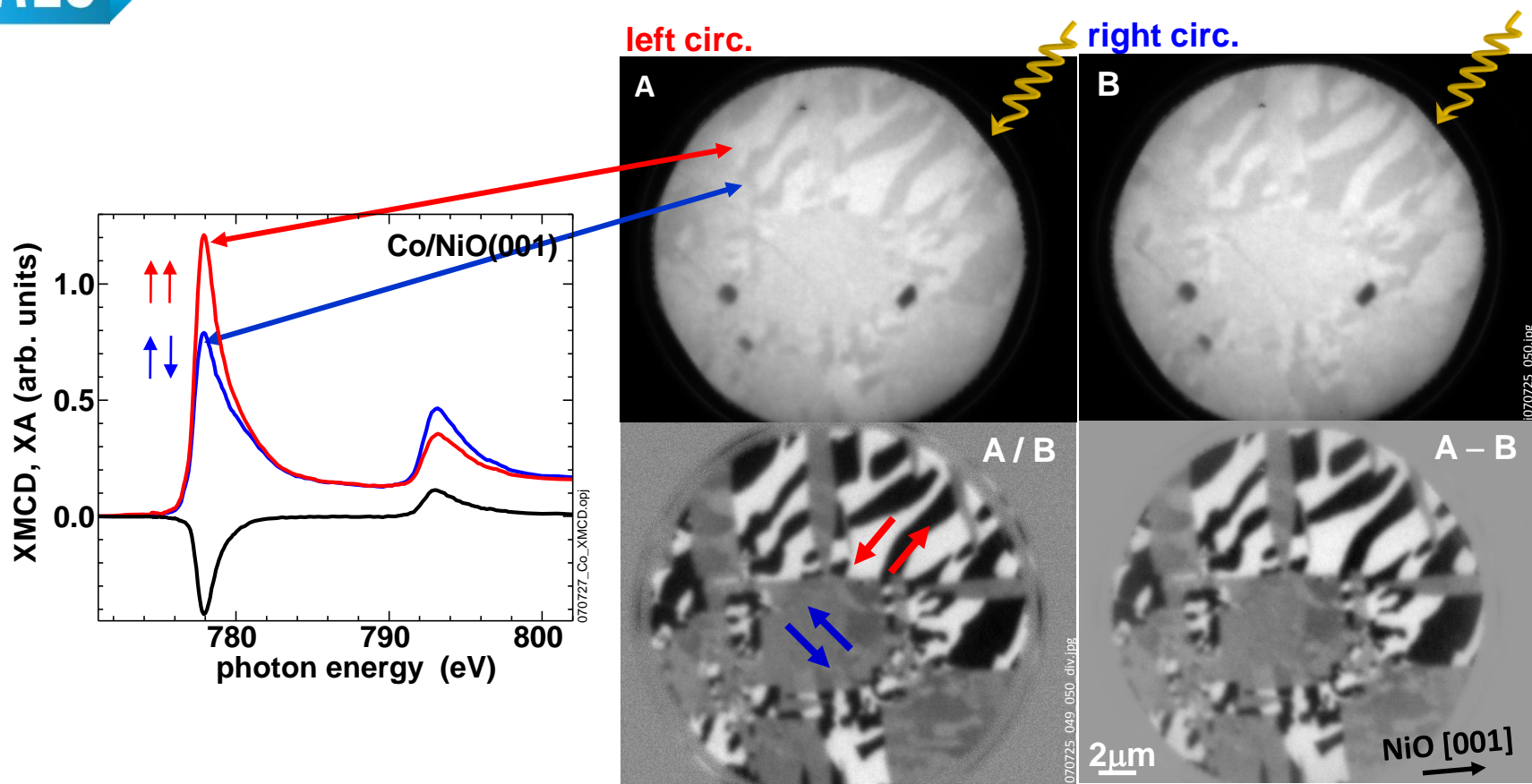
J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)





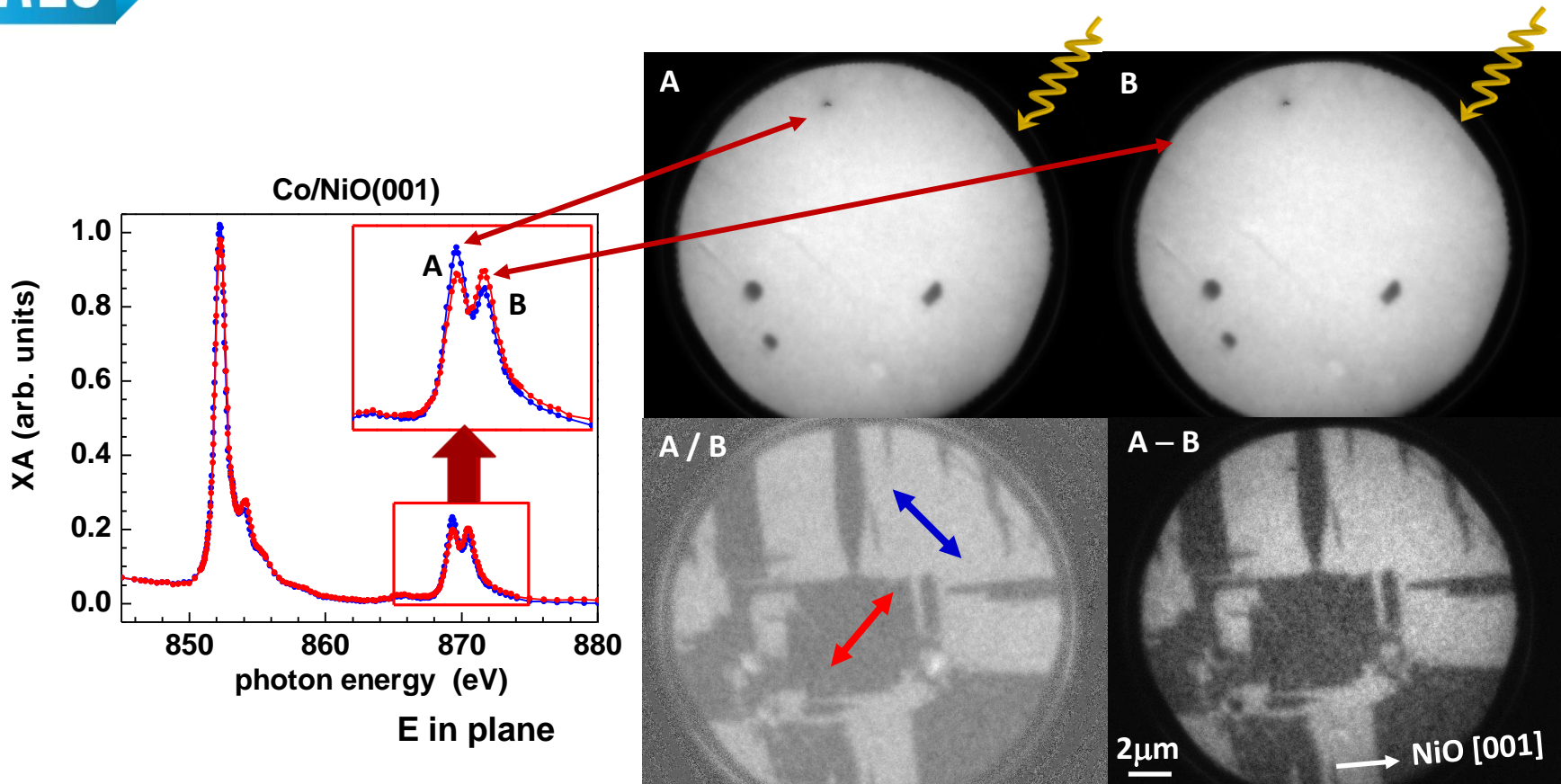
**10-50 nm spatial resolution**

**J. Stöhr, H.C. Siegmann,  
Magnetism (Springer)**



E. Arenholz *et al.*,  
*Appl. Phys. Lett.* **93**, 162506 (2008)

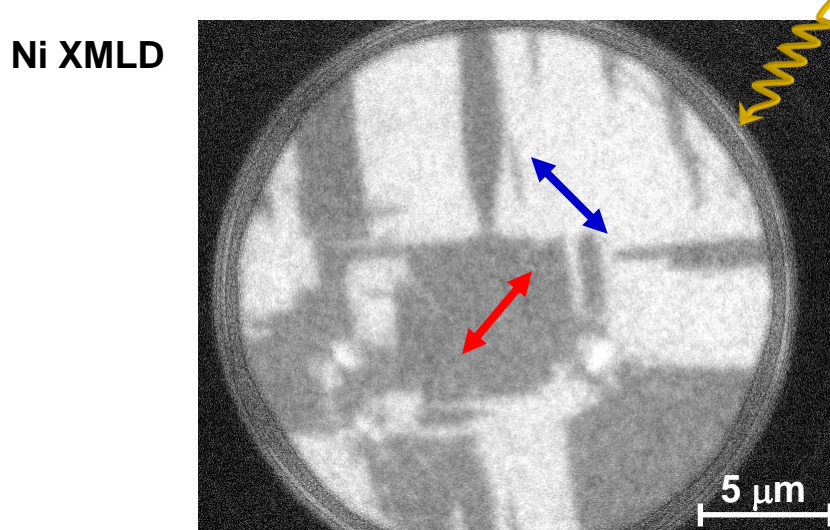
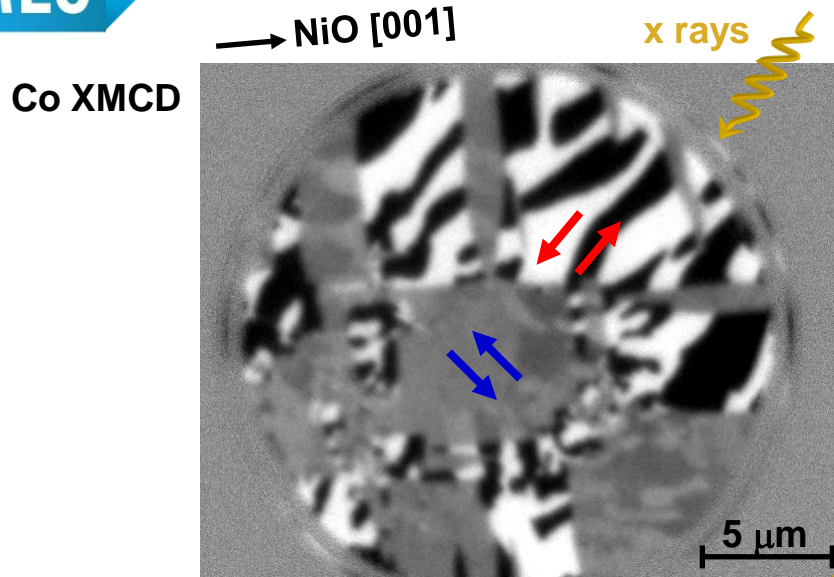
- + Images taken with left and right circularly polarized x-rays at photon energies with XMCD, i.e. Co  $L_3$  edge, provide magnetic contrast and domain images.



E. Arenholz *et al.*,  
*Appl. Phys. Lett.* **93**, 162506 (2008)

- + Images taken with linearly polarized x-rays at photon energies with XMLD, i.e. Ni  $L_2$  edge, provide magnetic contrast and domain images.





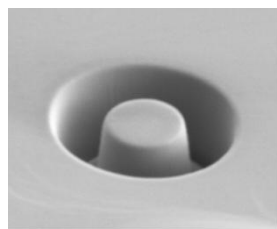
probing in-plane

- + Taking into account the geometry dependence of the Ni XMLD signal  
⇒ Perpendicular coupling of Co and NiO moments at the interface.

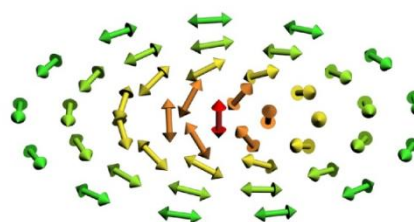
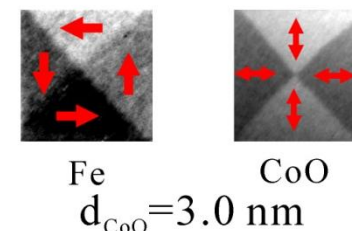
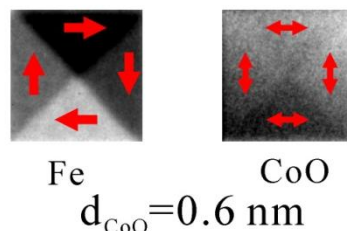
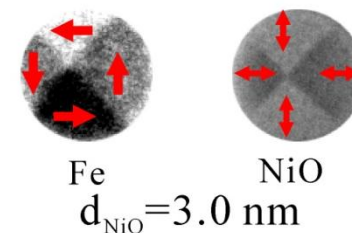
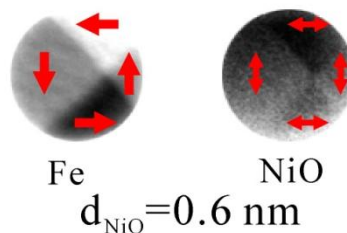
E. Arenholz *et al.*,  
Appl. Phys. Lett. 93, 162506 (2008)

+ First direct observation of vortex state in antiferromagnetic CoO and NiO disks in Fe/CoO and Fe/NiO bilayers using XMCD and XMLD.

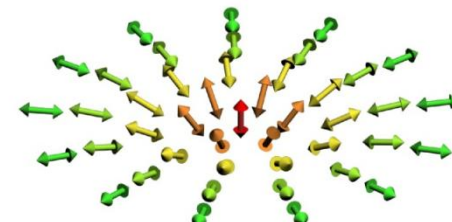
- + Two types of AFM vortices:
- conventional curling vortex as in ferromagnets
  - divergent vortex, forbidden in ferromagnets
  - thickness dependence of magnetic interface coupling



2  $\mu\text{m}$



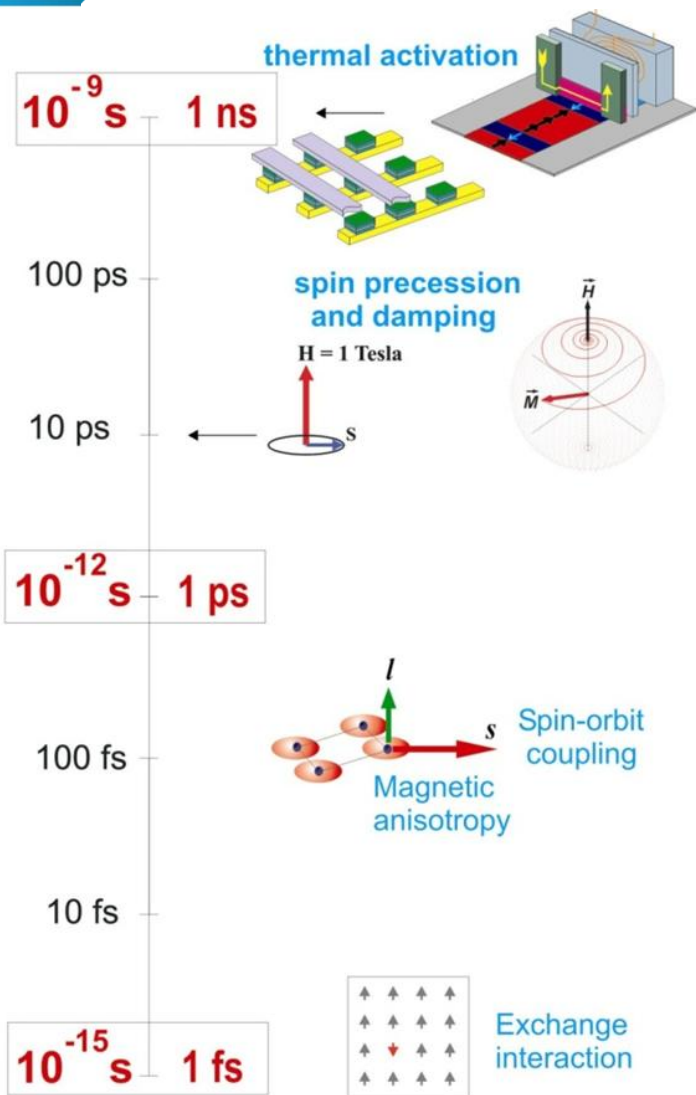
conventional curling vortex



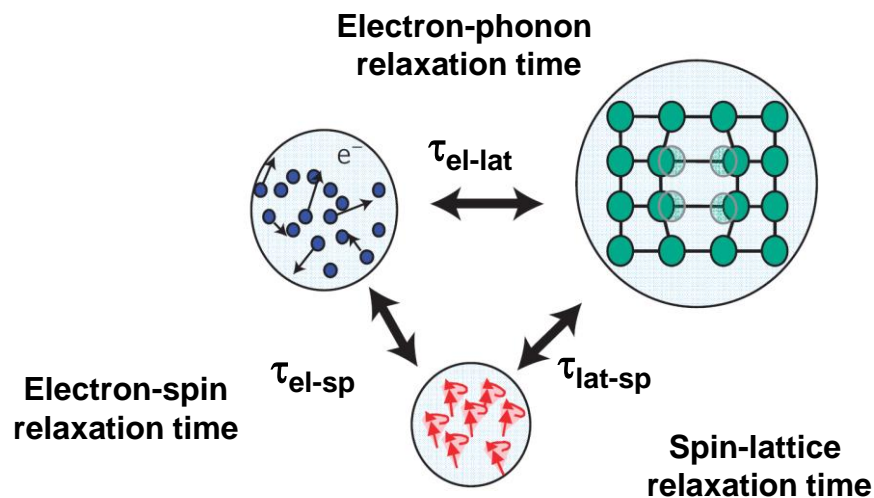
divergent vortex

J. Wu *et al.*,  
Nature Phys. **7**, 303 (2011)



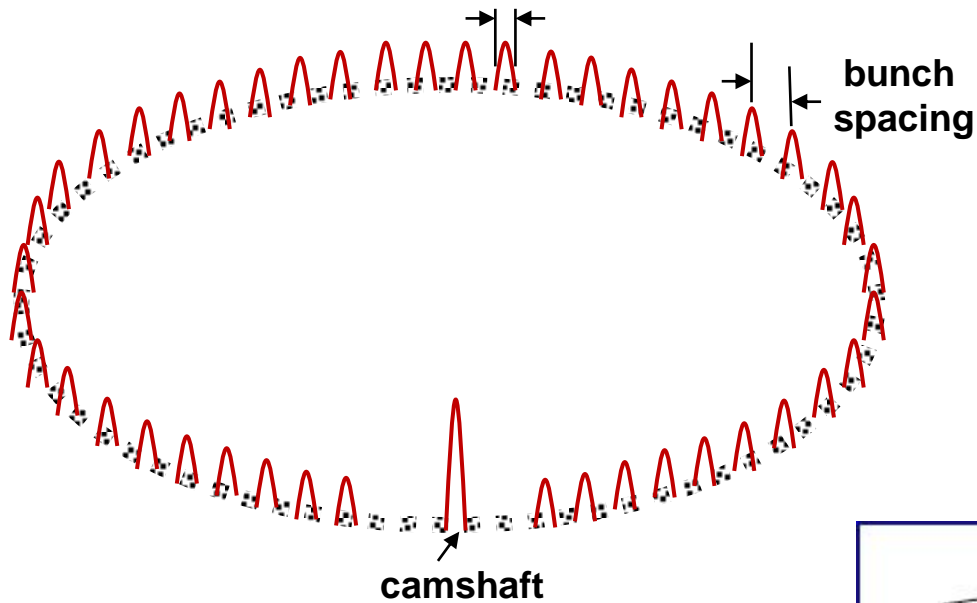


- + Energy reservoirs in a ferromagnetic metal
- + Deposition of energy in one reservoir
- ⇒ Non-equilibrium distribution and subsequent relation through energy and angular momentum exchange



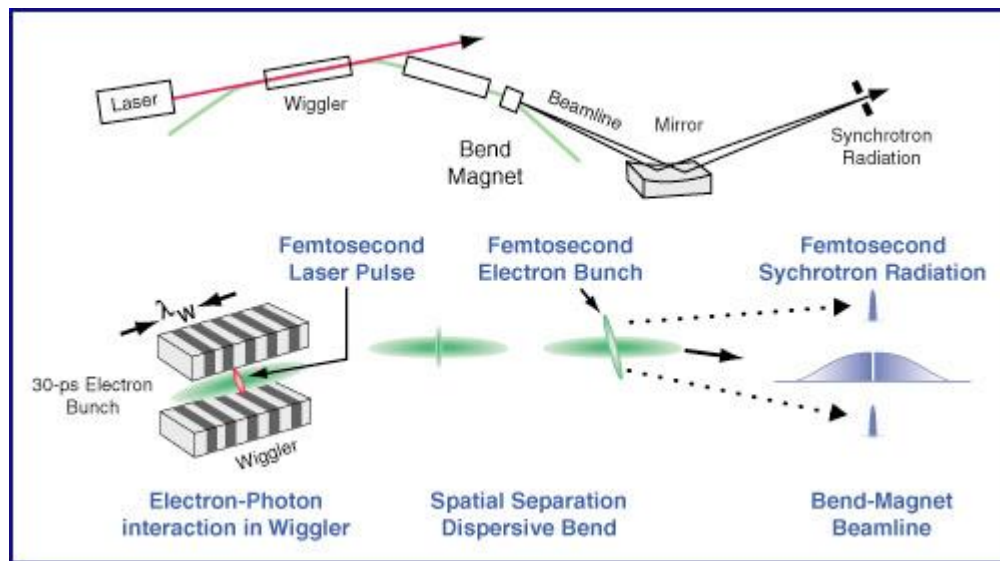
J. Stöhr, H.C. Siegmann, Magnetism (Springer)

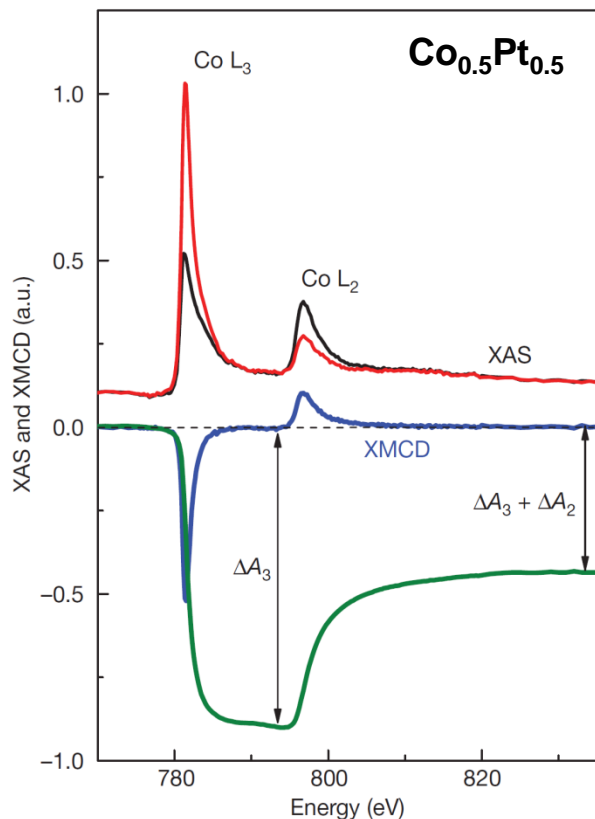
**Pulse length 70 ps**



- + 256-320 bunches for 500mA beam current
- + Possibility of one or two 5mA "camshaft" bunches in filling gaps
- + Bunch spacing:  
 multibunch mode: 2 ns  
 two-bunch mode: 328 ns
- + Pulse length 70ps

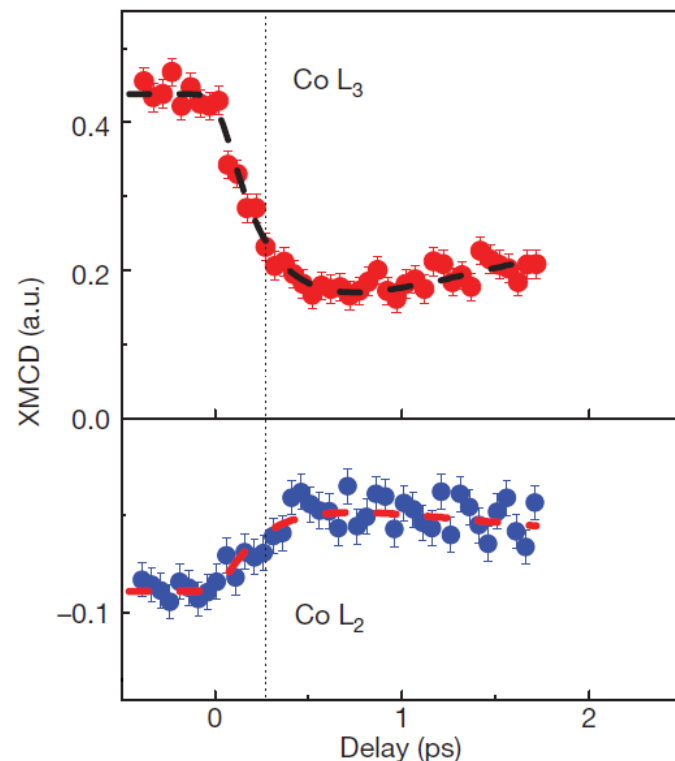
+ <300 fs x ray pulses though  
 "laser bunch-slicing technique"

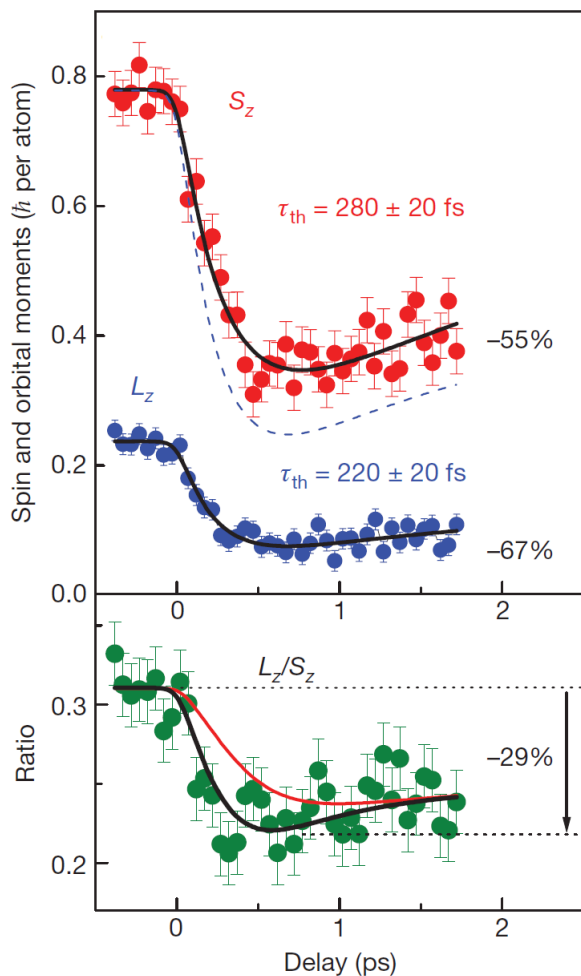




C. Boeglin, *et al.*,  
Nature **465**, 458 (2010)

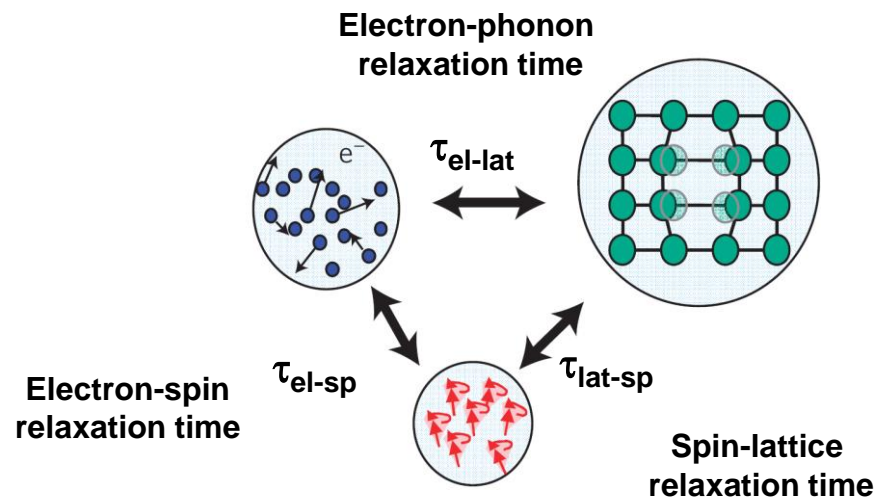
- + Orbital ( $L$ ) and spin ( $S$ ) magnetic moments can change with total angular momentum is conserved.
- + Efficient transfer between  $L$  and  $S$  through spin-orbit interaction in solids
- + Transfer between  $L$  and  $S$  occurs on fs timescales.
- +  $\text{Co}_{0.5}\text{Pt}_{0.5}$  with perpendicular magnetic anisotropy
- + 60 fs optical laser pulses change magnetization
- + Dynamics probed with XMCD using 120fs x-ray pulses
- + Linear relation connects  $\text{Co L}_3$  and  $\text{L}_2$  XMCD with  $L_z$  and  $S_z$  using sum rules

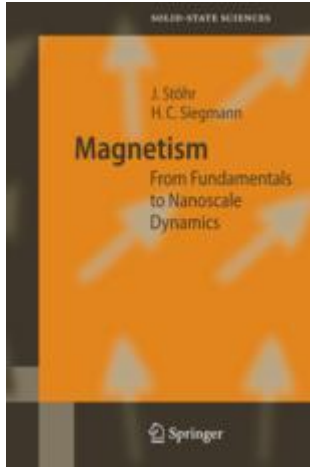




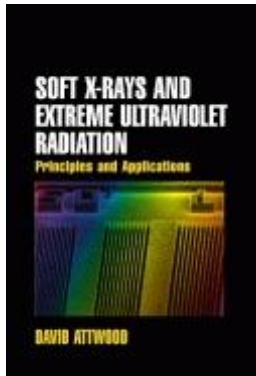
C. Boeglin, *et al.*,  
Nature **465**, 458 (2010)

- + Thermalization: Faster decrease of orbital moment
- + Theory: Orbital magnetic moment strongly correlated with magnetocrystalline anisotropy
- + Reduction in orbital moment  
 $\Leftrightarrow$  Reduction in magnetocrystalline anisotropy
- + Typically observed at elevated temperatures in static measurements as well





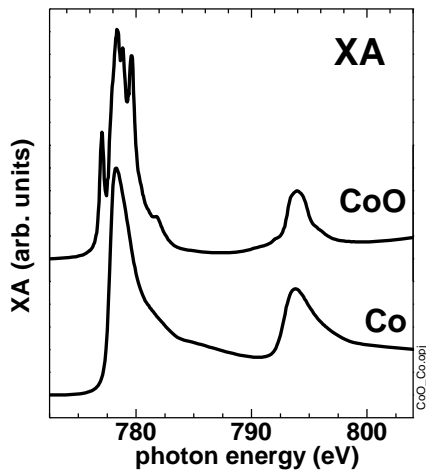
**J. Stöhr, H.C. Siegmann**  
**Magnetism– From Fundamentals to Nanoscale Dynamics**  
**Springer**



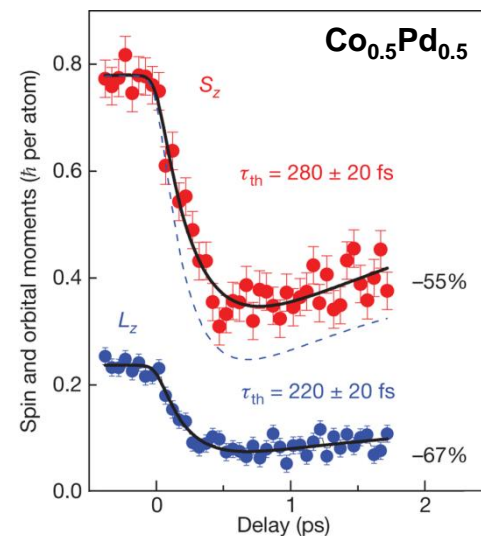
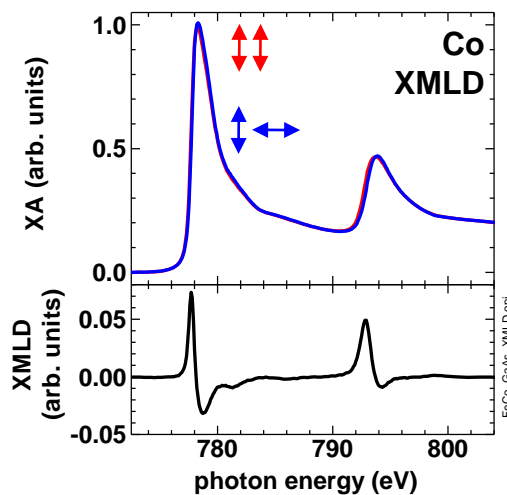
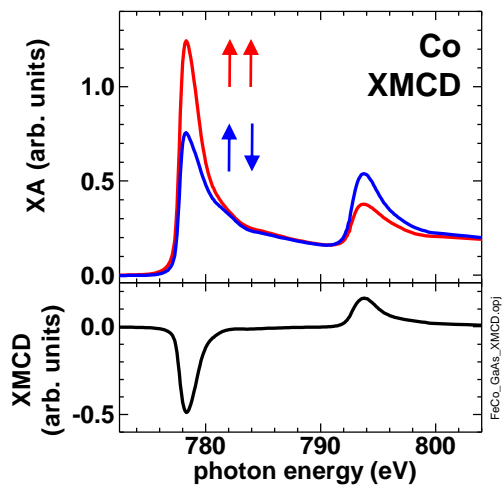
**D. Attwood**  
**Soft X-Rays and Extreme Ultraviolet Radiation:**  
**Principles and Applications**



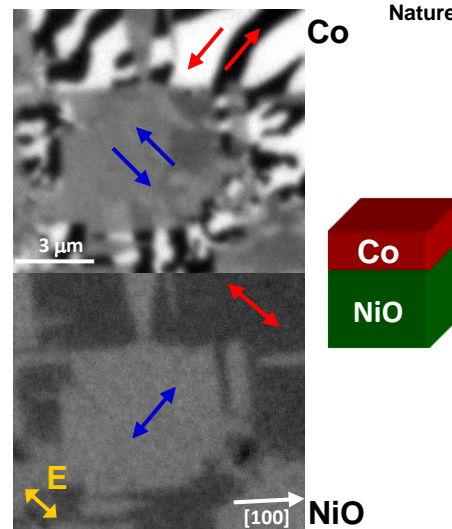
## Elke Arenholz, Advanced Light Source



- + X-ray absorption, XA
- + X-ray magnetic circular dichroism, XMCD
- + X-ray magnetic linear dichroism, XMLD
- + X-ray magnetic microscopy
- + Magnetization Dynamics



C. Boeglin *et al.*,  
Nature **465**, 458 (2011)



E. Arenholz *et al.*,  
Appl. Phys. Lett. **93**, 162506 (2008)